

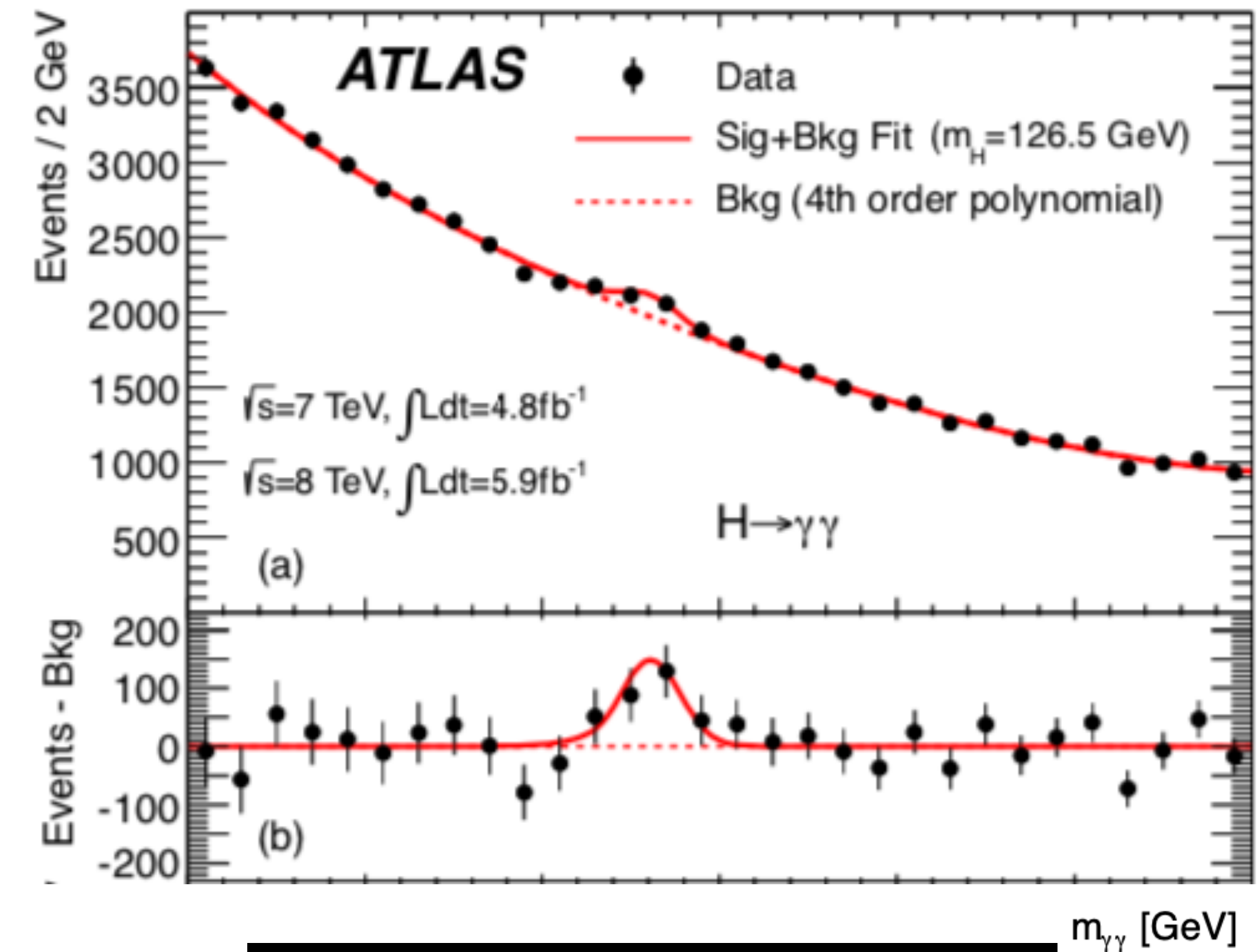
Studying the CP nature of the Top-Yukawa coupling

Federica Pasquali

MISP 2020

INTRODUCTION

- Higgs boson discovered in 2012

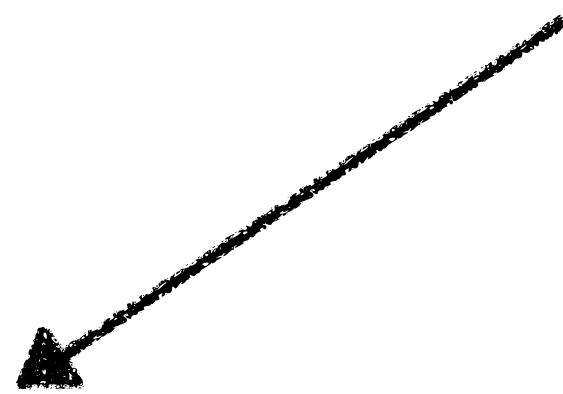


<https://arxiv.org/abs/1207.7214>

INTRODUCTION

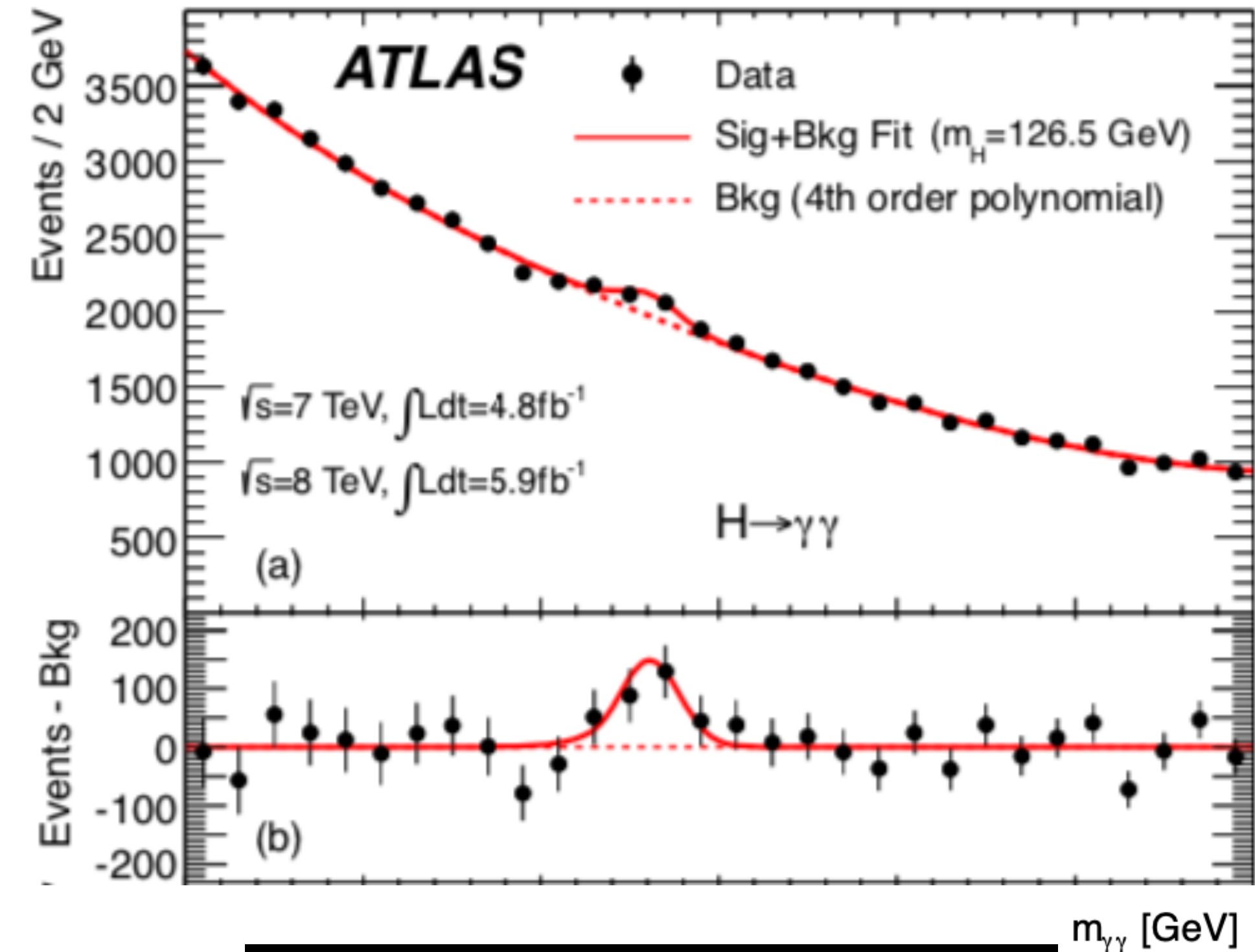
- Higgs boson discovered in 2012
- CP nature of the Higgs boson not yet well measured

- $\mathcal{L}_{Higgs} = |D_\mu \Phi|^2 - \mu^2 |\Phi|^2 - \lambda^4 |\Phi|^4$



Different options available

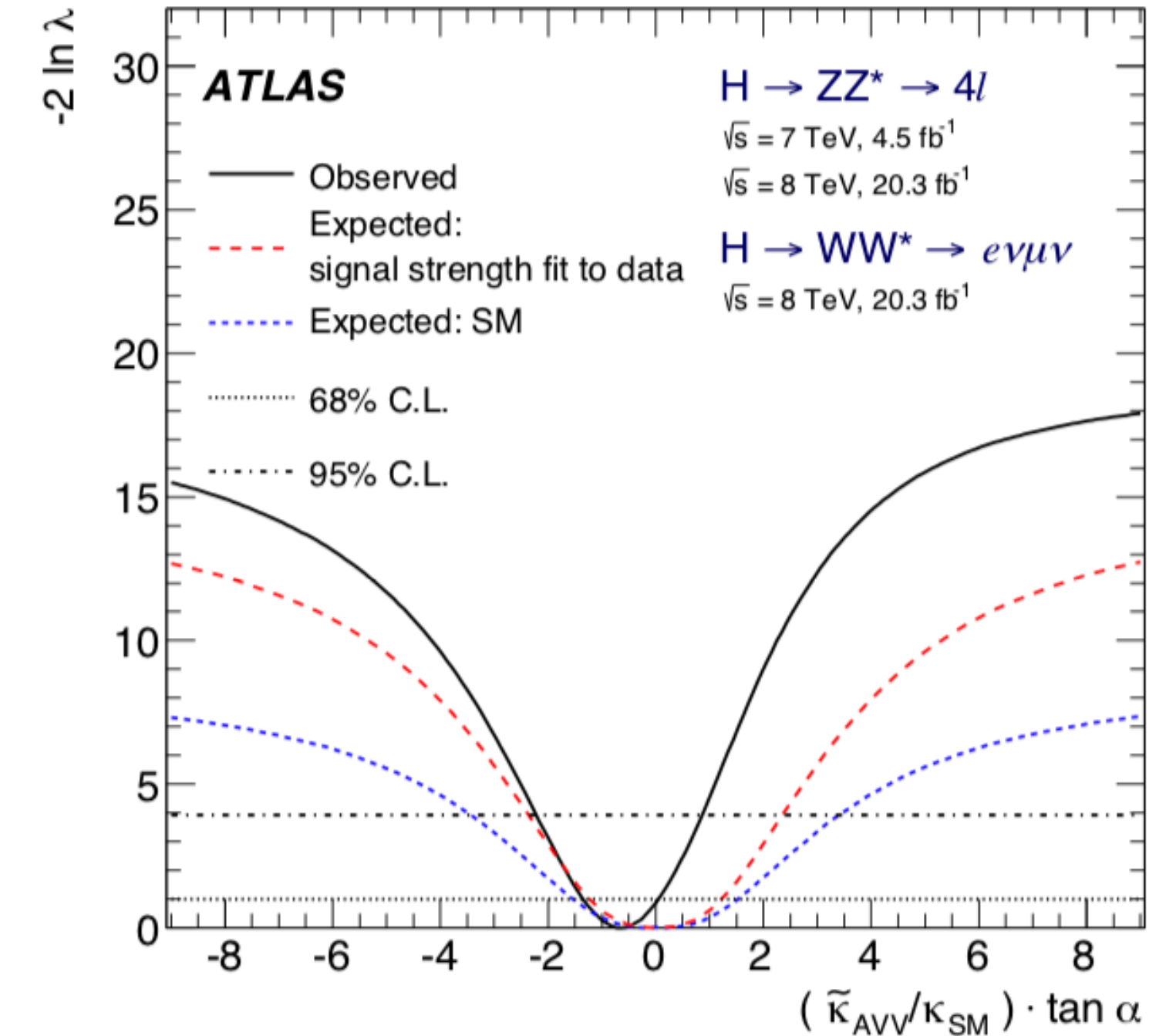
- CP even (SM) $CP |\Phi\rangle = |\Phi\rangle$
- CP odd $CP |\Phi\rangle = -|\Phi\rangle$
- CP mixed $CP |\Phi\rangle \neq \pm |\Phi\rangle$



<https://arxiv.org/abs/1207.7214>

THE AIM: SEARCH FOR CP VIOLATION

- Probe the CP nature of the Higgs boson in a model-independent way
- A non CP even Higgs boson hints at BSM physics
- In Run1 (0 jet channel)
(combination of WW^* , ZZ^* , $\gamma\gamma$)
 - Excluded the pure CP odd scenario
 - Up to 30% “CP odd” possible in mixed scenario (95%CL)



<https://arxiv.org/pdf/1506.05669.pdf>

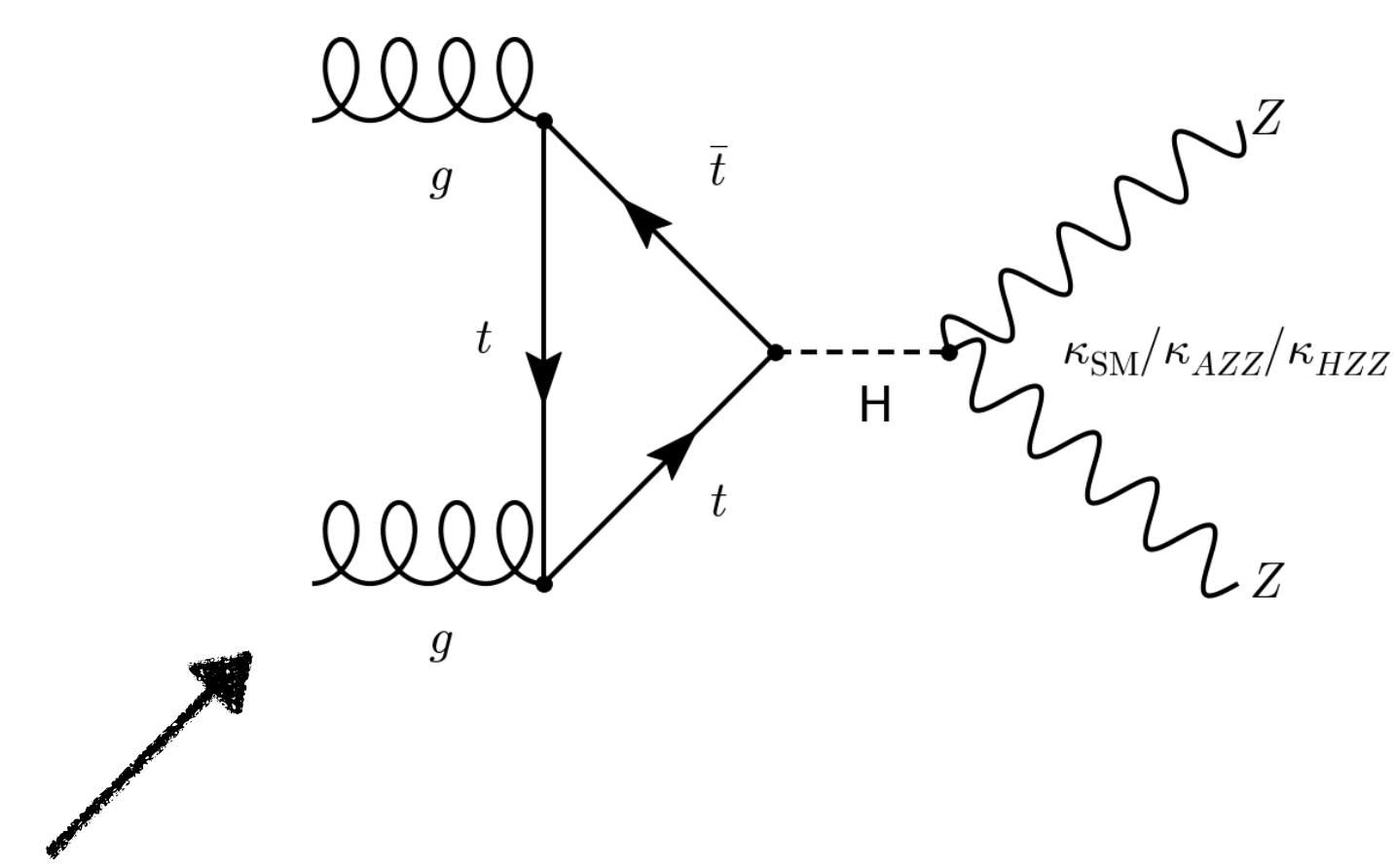
- In Run2 many CP analyses in different channels
 - $H \rightarrow ZZ^* \rightarrow 4l$, published with 36/fb <https://arxiv.org/pdf/1712.02304.pdf>



in this talk

$H \rightarrow ZZ^* \rightarrow 4l$ - MAIN PICTURE

- Why?
 - clear signature and high signal-to-background ratio
- Analysis strategy
 - events are classified according their different production modes (ggF, VBF, VH)
 - assuming that the BSM changes in the Higgs sector do not to affect the SM background processes
- Signal samples
 - the signal is modelled by using POWHEG Monte Carlo at NLO with QCD corrections up to NLO and QCD soft-gluon re-summations up NNL logarithm applied
 - morphing technique used to interpolate the BSM signal scenarios
- Main Background
 - continuum $ZZ^{(*)}$ production
 - Use of BDT to better separate signal and background
- A maximum likelihood (ML) fit is used for the statistical interpretation of the results



$H \rightarrow ZZ^* \rightarrow 4l$ - MAIN PICTURE

- Event Selection

- 4 leptons selected

- $4\mu, 2e2\mu, 2\mu2e, 4e$

- The three highest- p_T leptons required to have $p_T > 20\text{GeV}$, $p_T > 15\text{GeV}$, $p_T > 10\text{GeV}$, respectively

- The leading lepton pair required to have $50 < m_{1,2} < 106 \text{ GeV}$

- The sub-leading lepton pair $m_{min} < m_{3,4} < 115 \text{ GeV}$

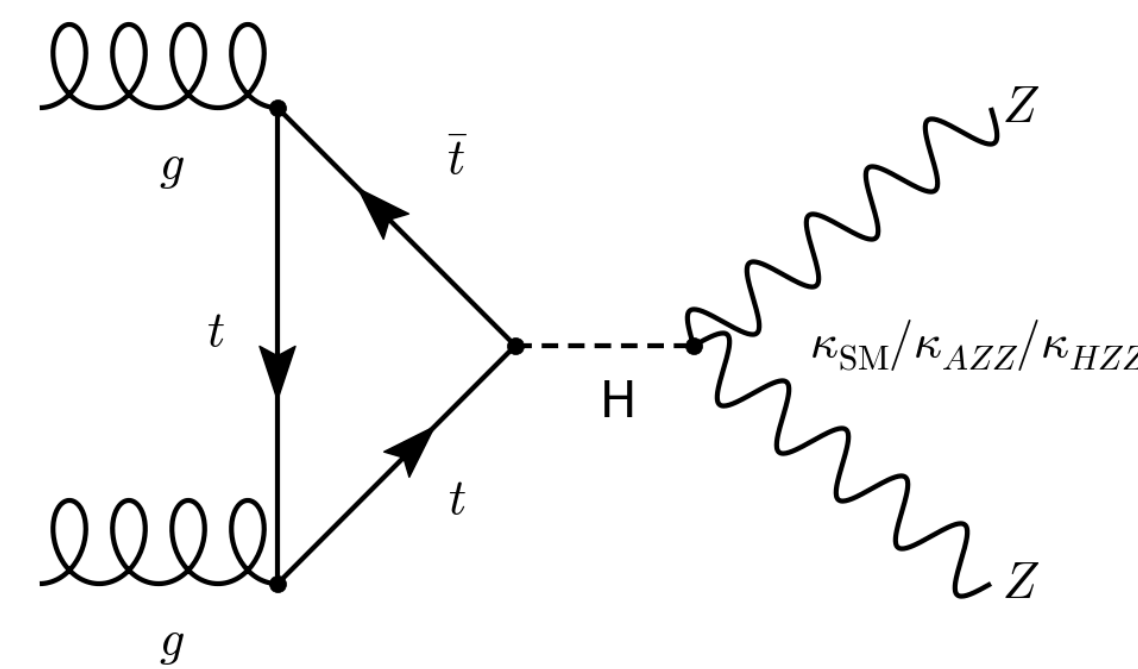
- Where m_{min} goes from 12 to 50 GeV depending on the four-lepton invariant mass m_{4l}

- The two-lepton pairs must have an angular separation of $\Delta R = \sqrt{(\Delta y)^2 + (\Delta\phi)^2} > 0.1(0.2)$

for same-flavour (different-flavour) lepton pairs

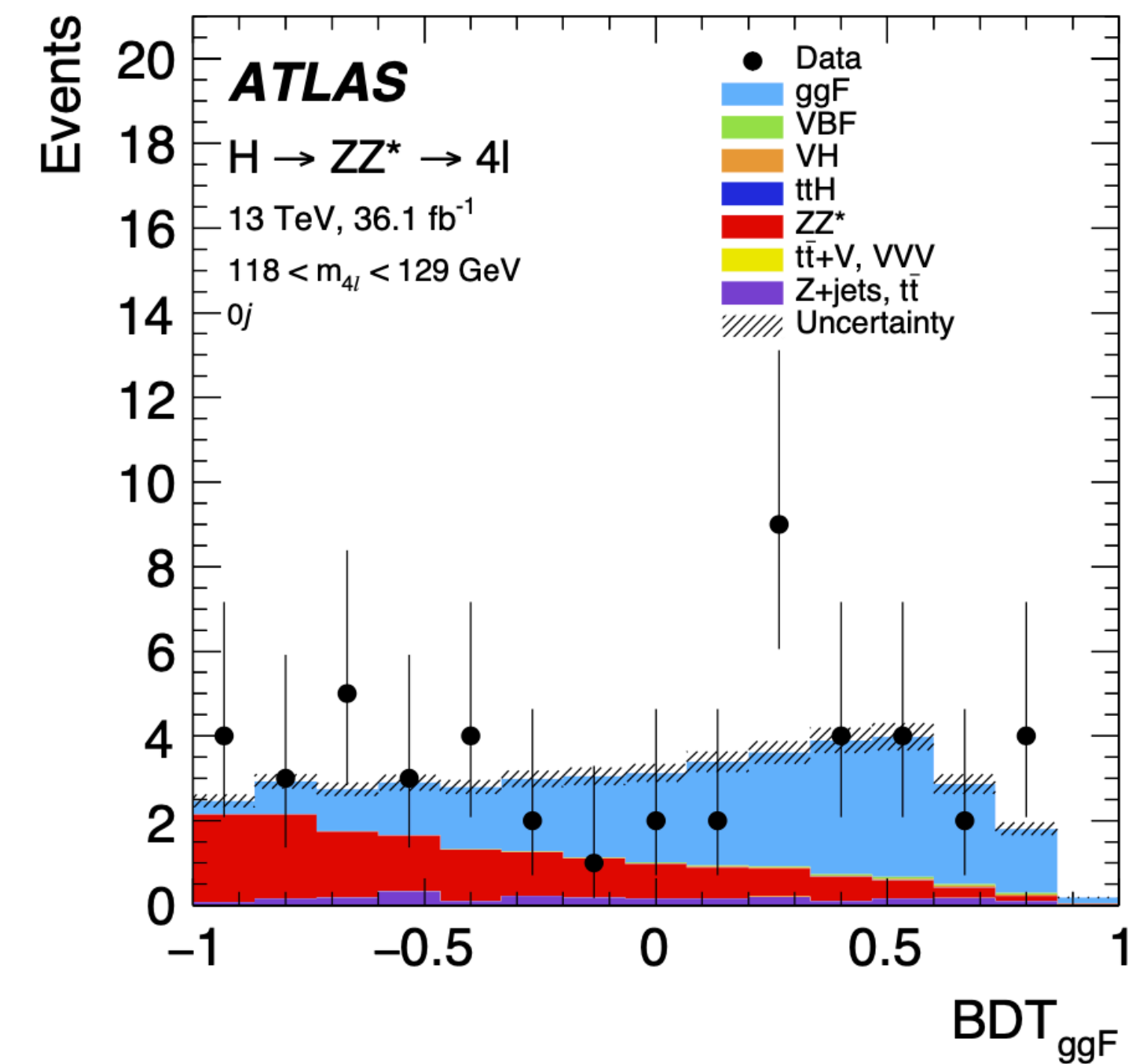
- Each electron (muon) must have a transverse impact parameter significance

$$|d_0|/\sigma(d_0) < 5(3)$$



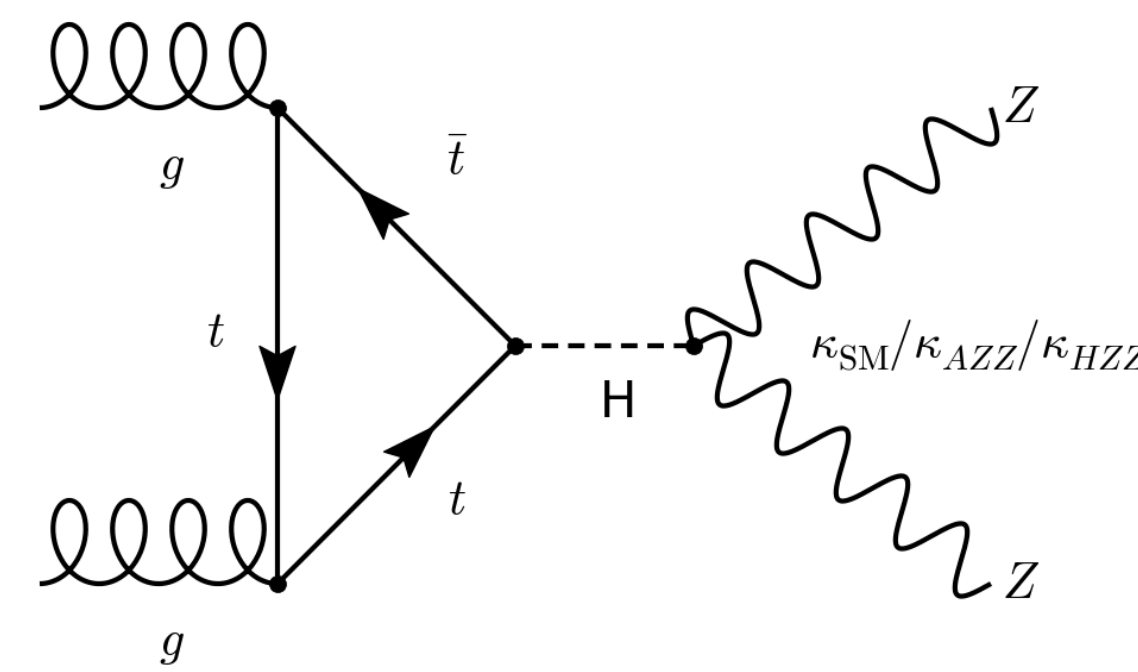
H \rightarrow ZZ* \rightarrow 4l - BDT

- Differences between H \rightarrow ZZ* \rightarrow 4l and the ZZ* background incorporated in a BDT discriminant
 - Training done with fully simulated
 - signal events with $m_H = 125$ GeV
 - $qq \rightarrow ZZ^*$ continuous background distribution
 - Only events with $118 < m_{4l} < 129$ GeV
 - asymmetric around 125 GeV to include the residual effects of bremsstrahlung
 - Variables used
 - p_T^{4l} : transverse momentum of the four-lepton system
 - η_T^{4l} : pseudorapidity of the four-lepton system
 - D_{ZZ^*} defined as $\ln(|\mathcal{M}_{sig}|^2 / |\mathcal{M}_{ZZ}|^2)$ with \mathcal{M}_x the matrix element for the considered process



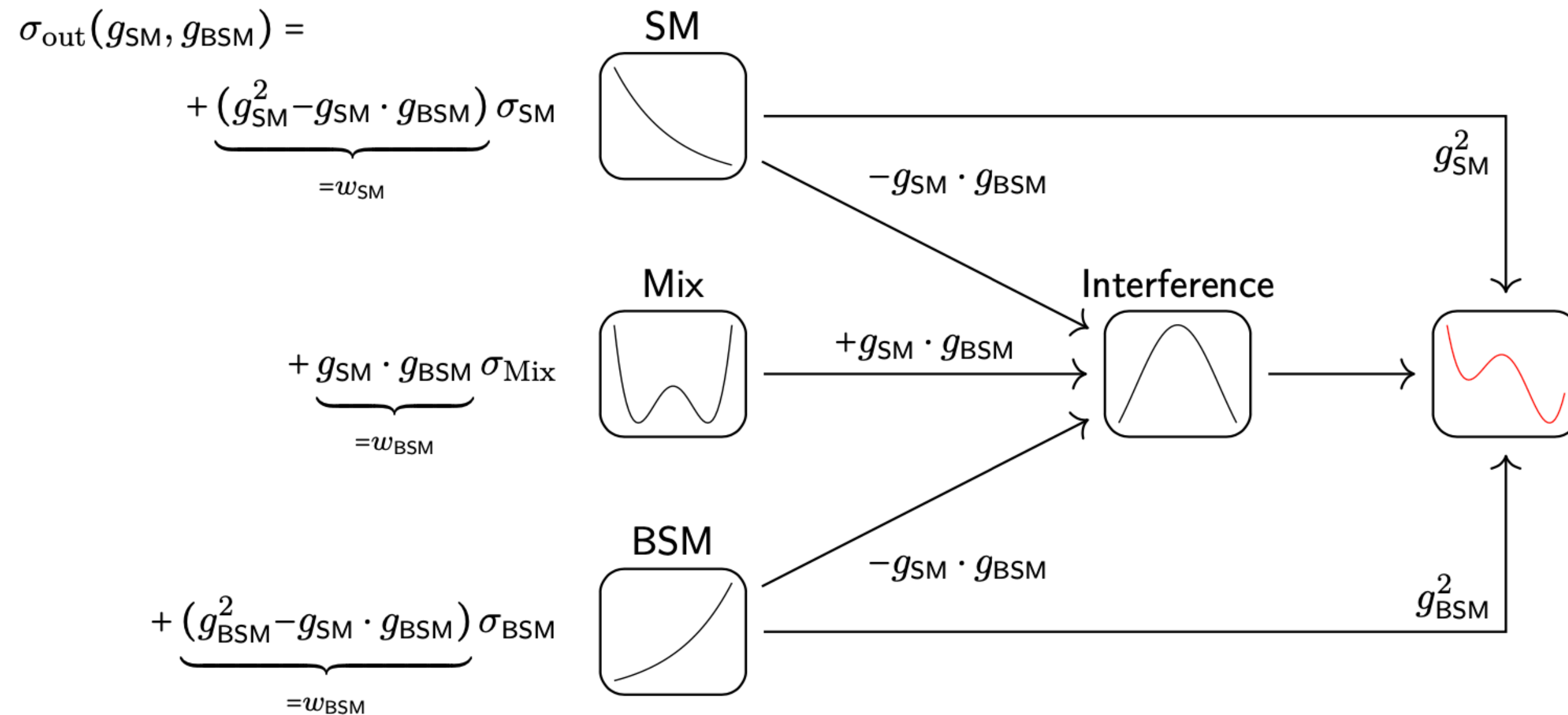
H \rightarrow ZZ* \rightarrow 4l - MORPHING

- The morphing functionality allows to interpolate/extrapolate a given physical (possibly multi-dimensional) observable as a function of SM-like and BSM Higgs boson couplings to known particles.



<http://cdsweb.cern.ch/record/2066980>

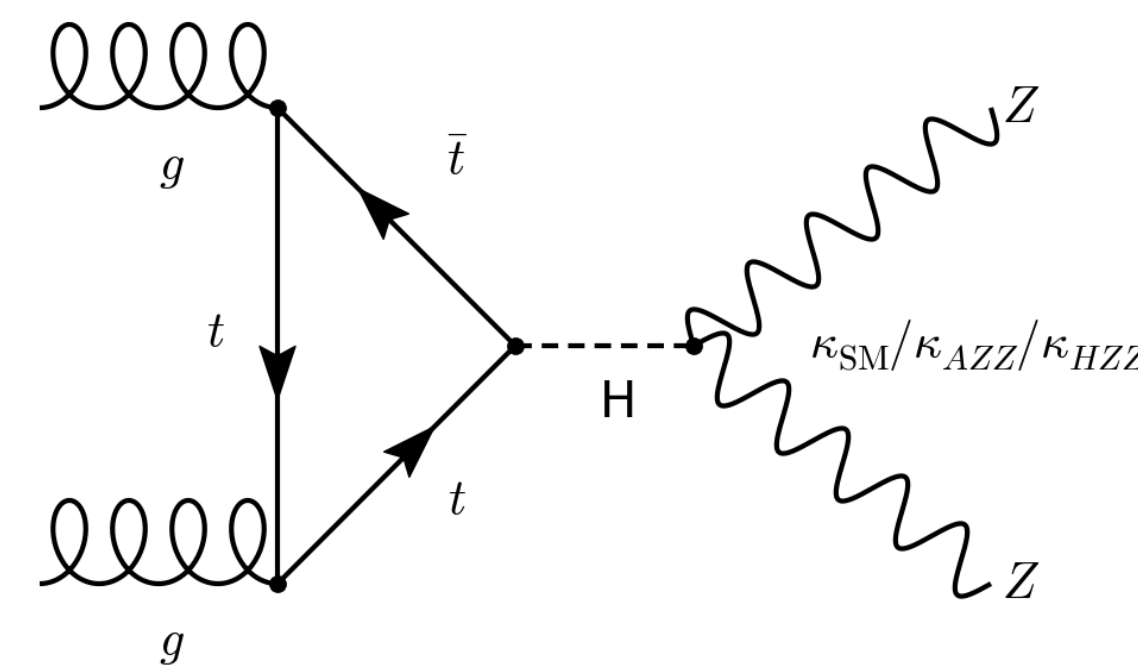
$$\begin{array}{c|cc} & g_{SM} & g_{BSM} \\ \hline SM & 1 & 0 \\ Mix & 1 & 1 \\ BSM & 0 & 1 \end{array} \xrightarrow{\text{polynomials}} \begin{array}{c|ccc} & g_{SM}^2 & g_{SM}g_{BSM} & g_{BSM}^2 \\ \hline SM & 1 & 0 & 0 \\ Mix & 1 & 1 & 1 \\ BSM & 0 & 0 & 1 \end{array} \xrightarrow{\text{inversion}} \begin{array}{c|ccc} & SM & Mix & BSM \\ \hline g_{SM}^2 & 1 & 0 & 0 \\ g_{SM}g_{BSM} & -1 & 1 & -1 \\ g_{BSM}^2 & 0 & 0 & 1 \end{array}$$



N of base sample	ggF
$C_\alpha K_{H\nu\nu}$	3
$S_\alpha K_{A\nu\nu}$	3
$S_\alpha K_{A\nu g}$	3
$C_\alpha K_{H\nu\nu}$ VS. $S_\alpha K_{A\nu\nu}$	6
$C_\alpha K_{H\nu\nu}$ VS. $S_\alpha K_{A\nu g}$	9

H \rightarrow ZZ* \rightarrow 4l - HOW

- Measuring the CP state of the top-Yukawa coupling by using an EFT (effective field theory) approach



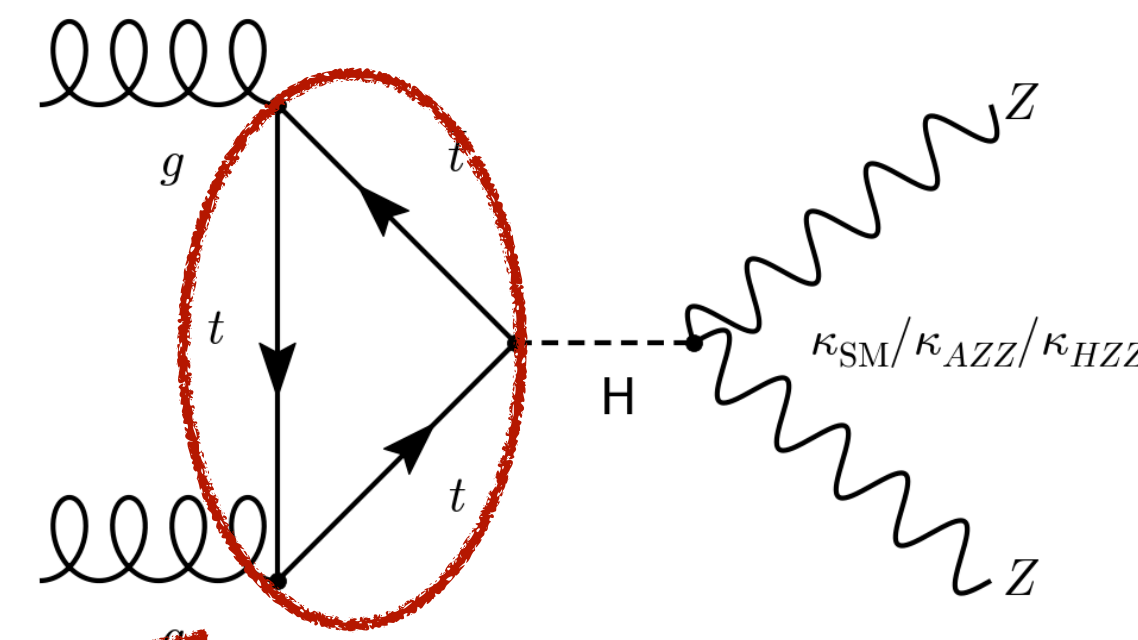
Higgs Caraterization model
<https://arxiv.org/abs/1306.6464>

$$\mathcal{L}_0^V = \left\{ \begin{aligned} &\kappa_{\text{SM}} \left[\frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \\ &- \frac{1}{4} \left[\kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + \tan \alpha \kappa_{A_{gg}} g_{A_{gg}} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right] \\ &- \frac{1}{4\Lambda} \left[\kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + \tan \alpha \kappa_{A_{ZZ}} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] \\ &- \frac{1}{2\Lambda} \left[\kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + \tan \alpha \kappa_{A_{WW}} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right] \end{aligned} \right\} \chi_0$$

- Λ set to 1 TeV, no new BSM particles below the energy scale Λ
- χ_0 bosonic state of spin 0
- α the mixing angle between the 0^+ and 0^- CP states
- κ_{HVV} , κ_{AVV} and $\kappa_{A_{gg}}$ are considered as possible BSM admixtures to the corresponding SM interactions
- $\kappa_{HWW} = \kappa_{HZZ} \equiv \kappa_{HVV}$ and $\kappa_{AWW} = \kappa_{AZZ} \equiv \kappa_{AVV}$
- κ_{SM} and κ_{Hgg} are fixed to the SM value of 1 if not stated otherwise
- CP even $\alpha = 0$
- CP odd $\alpha = 1$
- CP mixed $0 < \alpha < 1$

H $\rightarrow ZZ^* \rightarrow 4l$ - HOW

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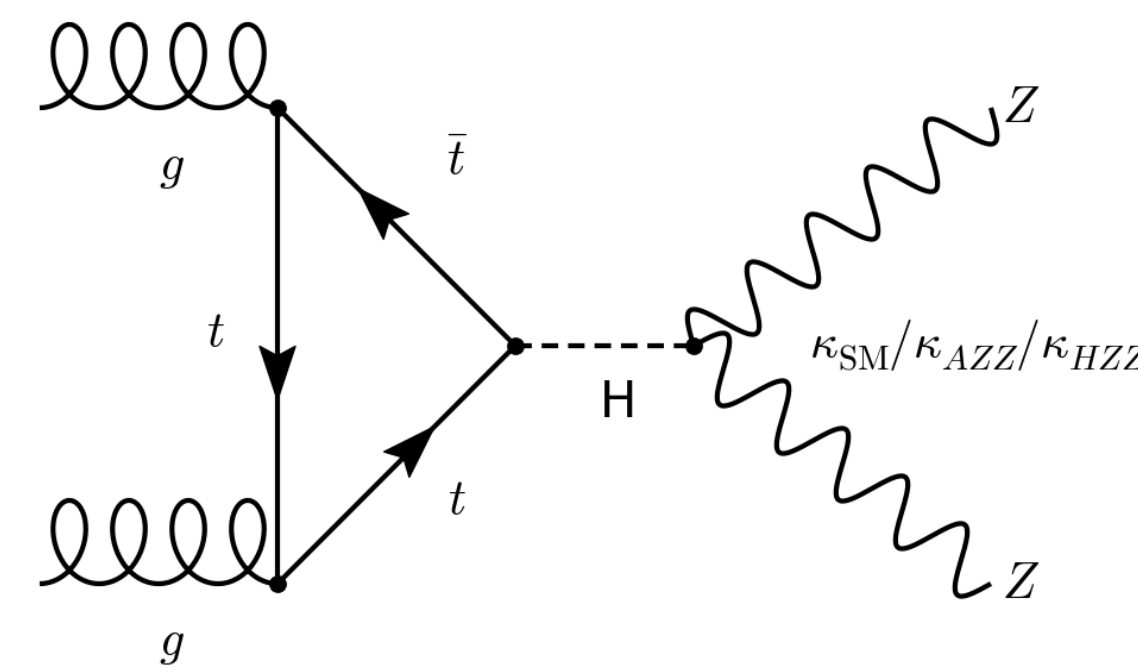
$$\mathcal{L}_0^V = \left\{ \kappa_{\text{SM}} \left[\frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right. \\
\left. - \frac{1}{4} \left[\kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + \tan \alpha \kappa_{A_{gg}} g_{A_{gg}} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right] \right. \\
\left. - \frac{1}{4} \frac{1}{\Lambda} \left[\kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + \tan \alpha \kappa_{A_{ZZ}} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] \right. \\
\left. - \frac{1}{2} \frac{1}{\Lambda} \left[\kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + \tan \alpha \kappa_{A_{WW}} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right] \right\}$$

This is what we are interested in

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- CP odd $\alpha = \pi$
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H \rightarrow ZZ* \rightarrow 4l - HOW

- Measuring the CP state of the top-Yukawa coupling by using an EFT (effective field theory) approach



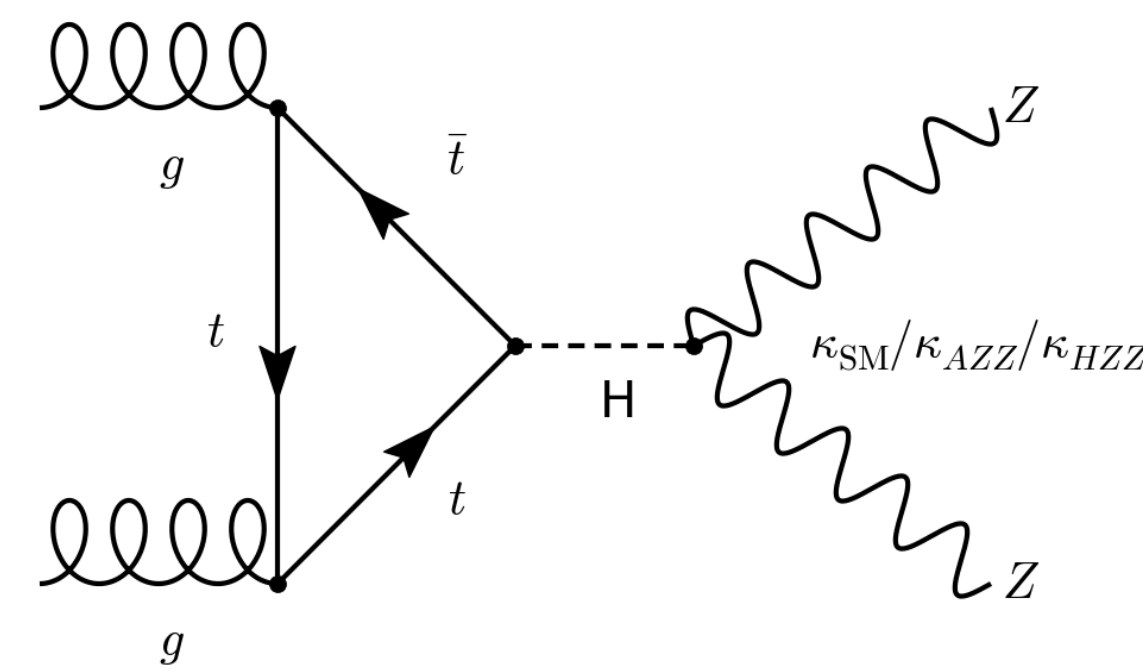
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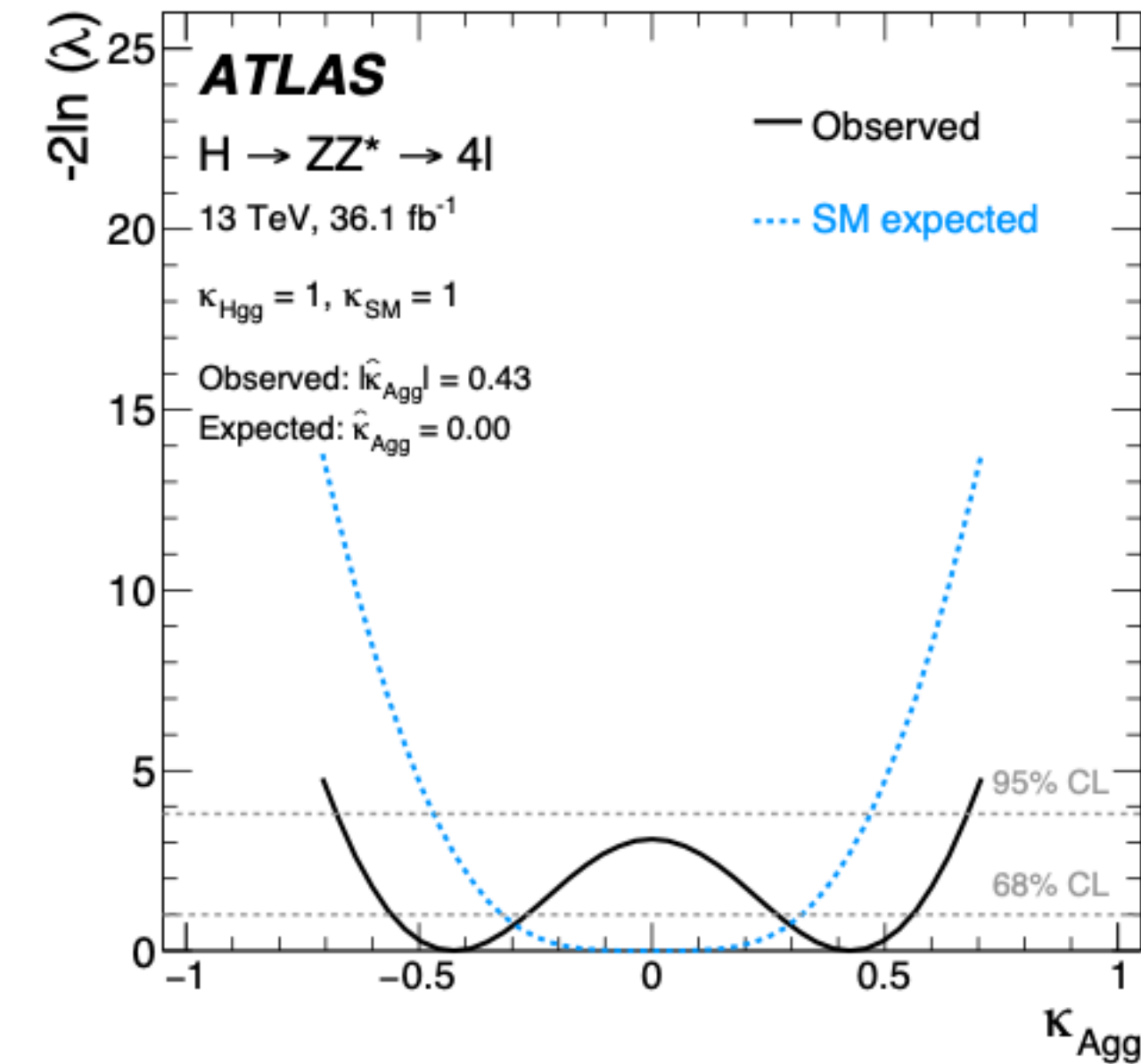
κ_{Agg} is measured assuming that all other BSM couplings are equal to 0

$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{Agg}$ FIT

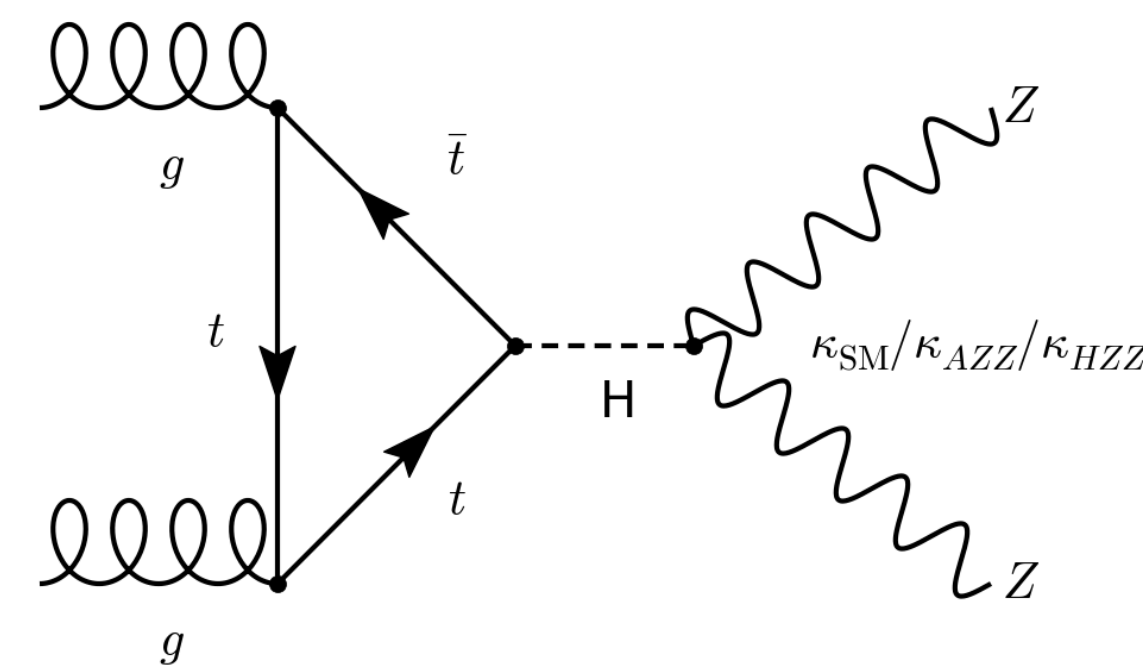


$$\mathcal{L}_0^V = -\frac{1}{4} [\kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + \tan\alpha \kappa_{Agg} A_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}]$$

- Probing BSM couplings in the ggH vertex
 - “Rate only” fit : obtained with an event yield-only analysis
 - The total width of the Higgs boson is modified to take into account the changes in the Branching Ratio due to the BSM contributions

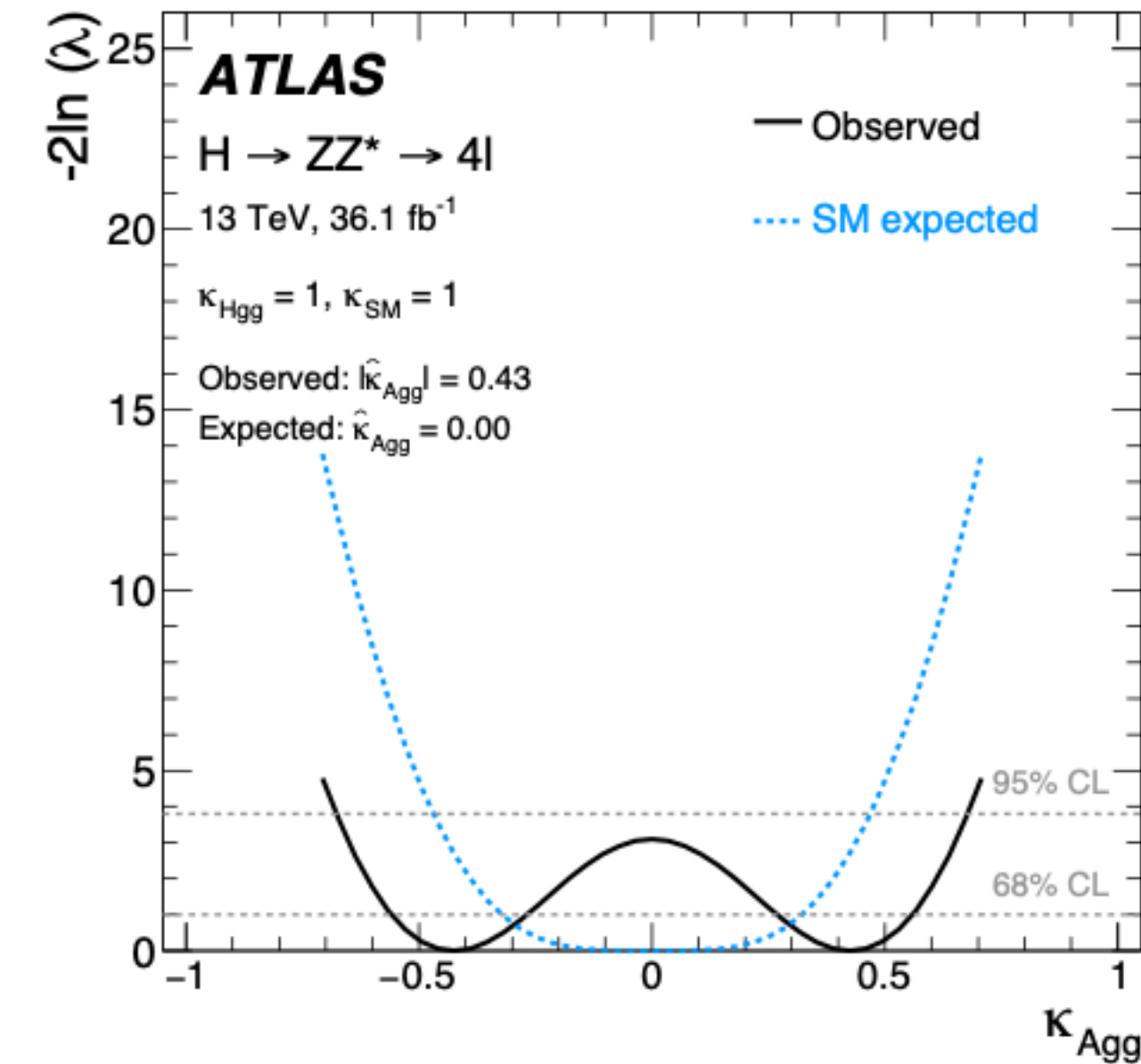


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$$\mathcal{L}_0^V = -\frac{1}{4} [\kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + \tan\alpha \kappa_{Agg} \tilde{A}_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}]$$

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BSM coupling	Fit configuration	Expected conf. inter.	Observed conf. inter.	Best-fit $\hat{\kappa}_{BSM}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
κ_{Agg}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	[-0.47, 0.47]	[-0.68, 0.68]	± 0.43	-	1.8σ

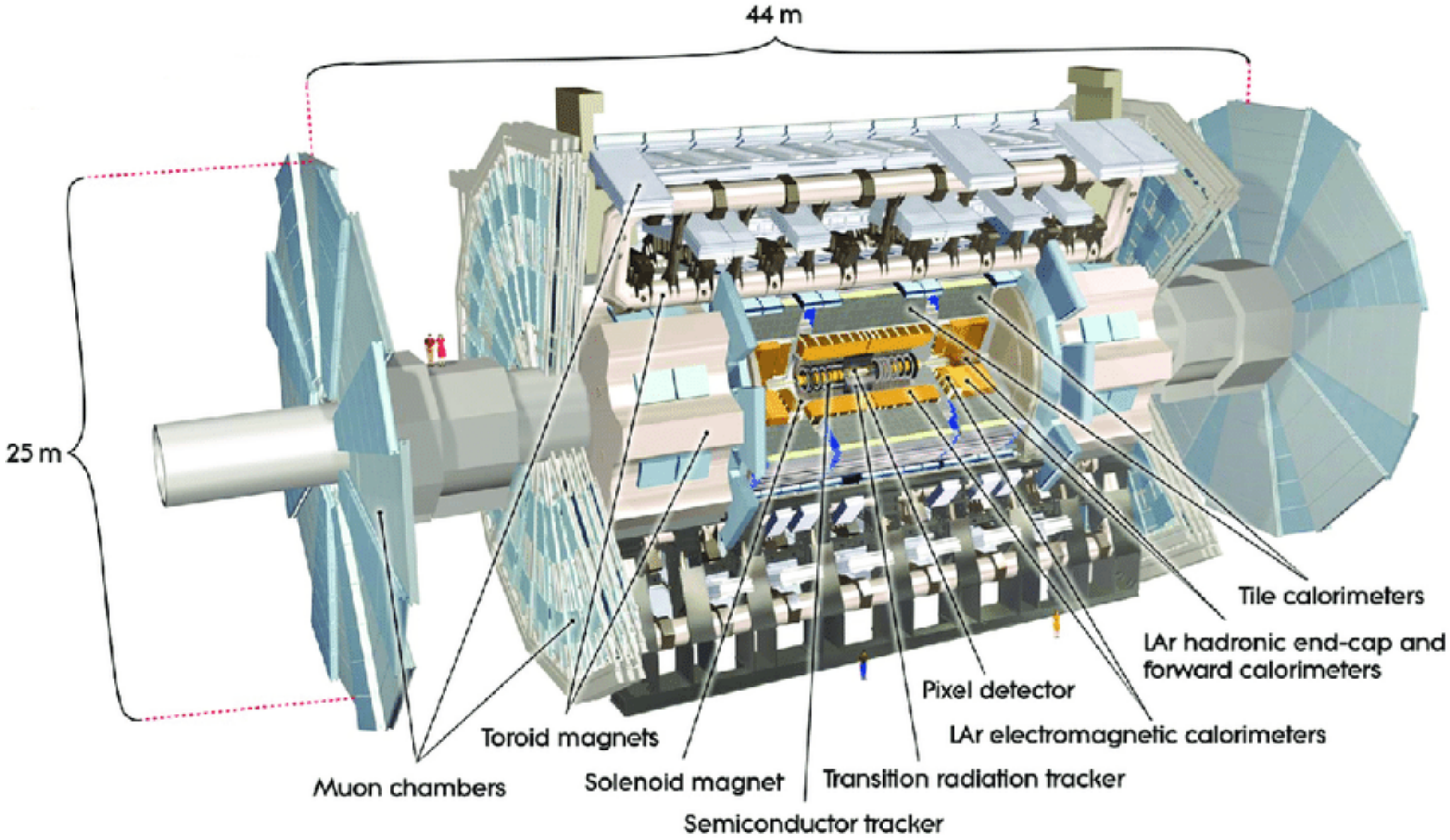
CONCLUSIONS

- Presented a study in Run2 of CP nature of the ggF production mode in $H \rightarrow ZZ^* \rightarrow 4l$
- Constraints are placed on possible BSM interactions of the Higgs boson within the framework of an effective Lagrangian extension of the SM
 - Pure CP odd scenario excluded
 - Up to 32% “CP odd” possible in mixed scenario (95%CL)
- Other measurements
 - $H \rightarrow \gamma\gamma$ <https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2019-029/>
 - $H \rightarrow WW^*$

Thanks!

Backup

ATLAS DETECTOR

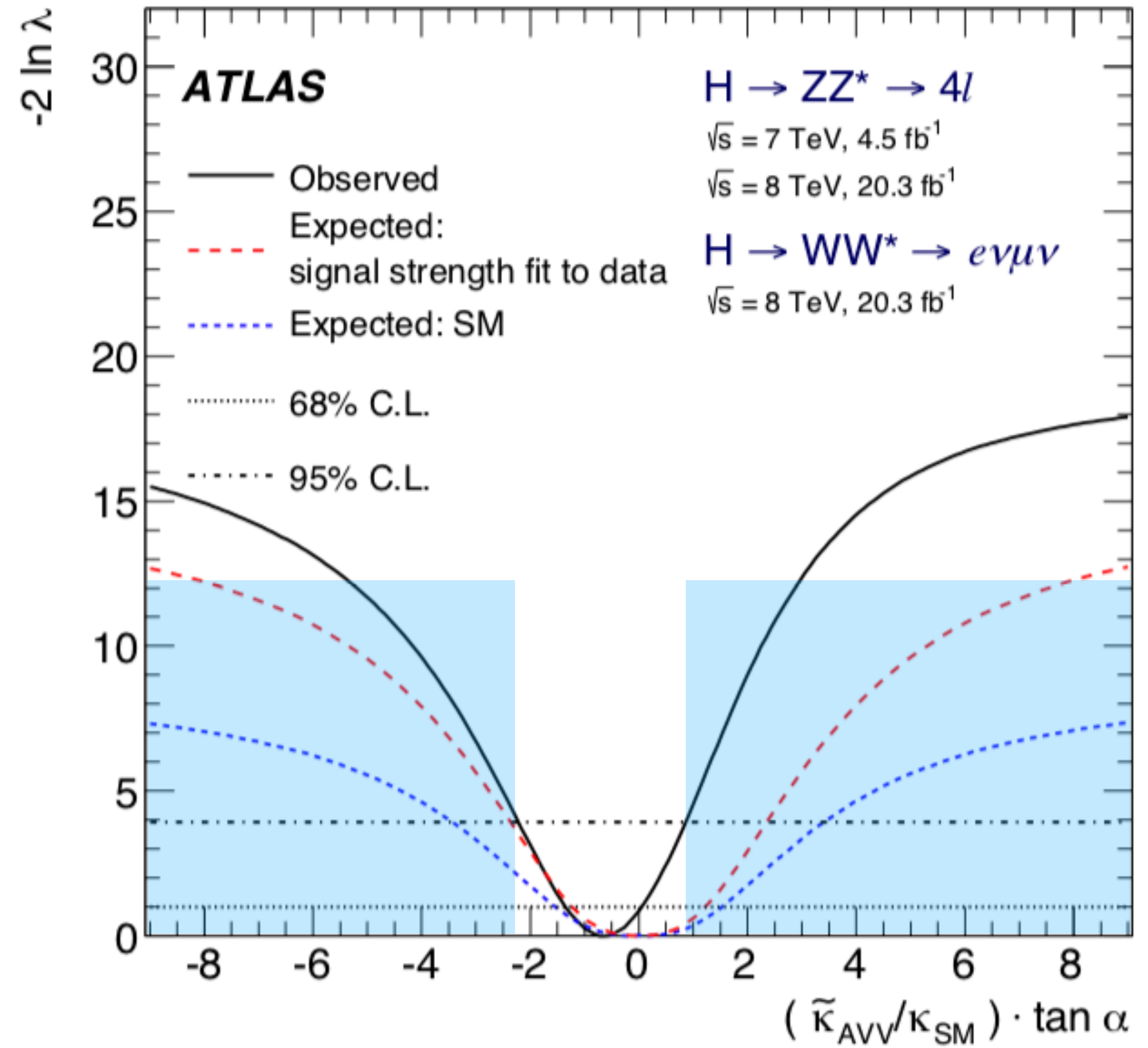


RUN1 EXCLUSION CALCULATION

$$\text{atan}(0.9) - \text{atan}(-2.3) = 1.89$$

$$\text{atan}(\text{inf}) - \text{atan}(-\text{inf}) = 6.28$$

$$1.89/6.28 \sim 30\%$$



TENSOR STRUCTURE

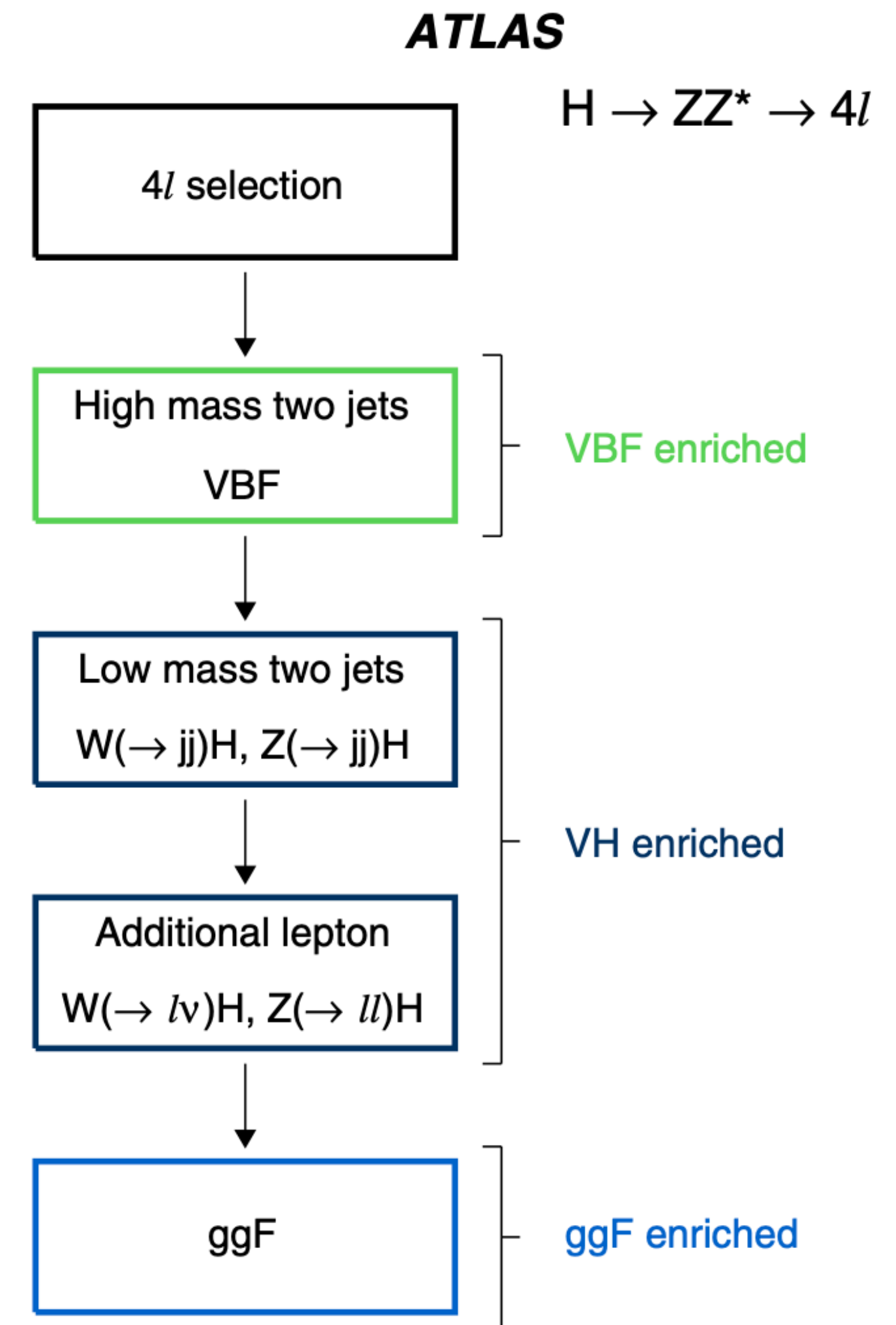
- V^μ represents the vector-boson field ($V = Z, W^\pm$), the $V^{\mu\nu}$ are the reduced field tensors and the dual tensor is defined as $\tilde{V}^{\mu\nu} = 1/2\epsilon^{\mu\nu\rho\sigma}V_{\rho\sigma}$
- $\kappa_{H\partial Z}$ and $\kappa_{H\partial W}$ corresponds to higher-order derivative operators, which have been neglected in this analysis
- Couplings are assumed to be real in order to ensure Lagrangian terms Hermitian
- $g_{HVV} \propto m_{Z/W}^2$

$$\mathcal{L}_0^V = \left\{ c_\alpha \kappa_{SM} \left[\frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] - \frac{1}{4} \frac{1}{\Lambda} \left[c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] - \frac{1}{2} \frac{1}{\Lambda} \left[c_\alpha \kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right] - \frac{1}{\Lambda} c_\alpha \left[\kappa_{H\partial Z} Z_\nu \partial_\mu Z^{\mu\nu} + \kappa_{H\partial W} (W_\nu^+ \partial_\mu W^{-\mu\nu} + h.c.) \right] \right\} \mathcal{X}_0$$

CLASSIFICATION OF THE HIGGS BOSON PRODUCTION MODES

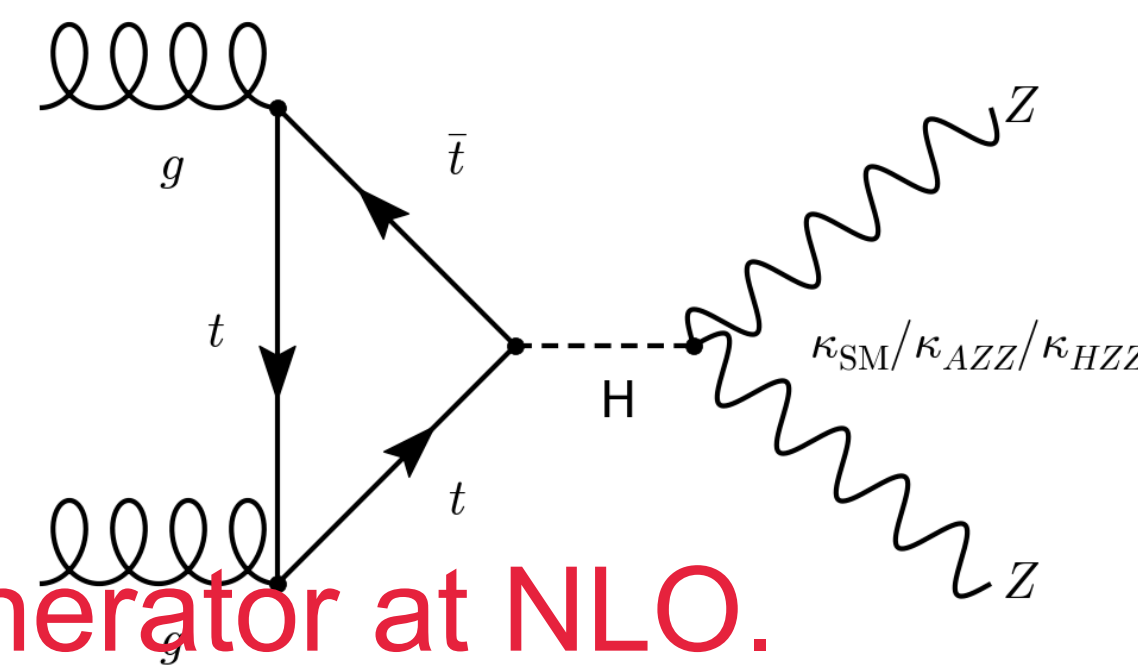
- In base of the production processes exclusive regions of the phase space are defined. These ones are called “production bins”.
- The bins are chosen in such a way that the measurement precision is maximized and at the same time possible BSM contributions can be isolated.

CATEGORIES



$H \rightarrow ZZ^* \rightarrow 4l$ - MAIN PICTURE

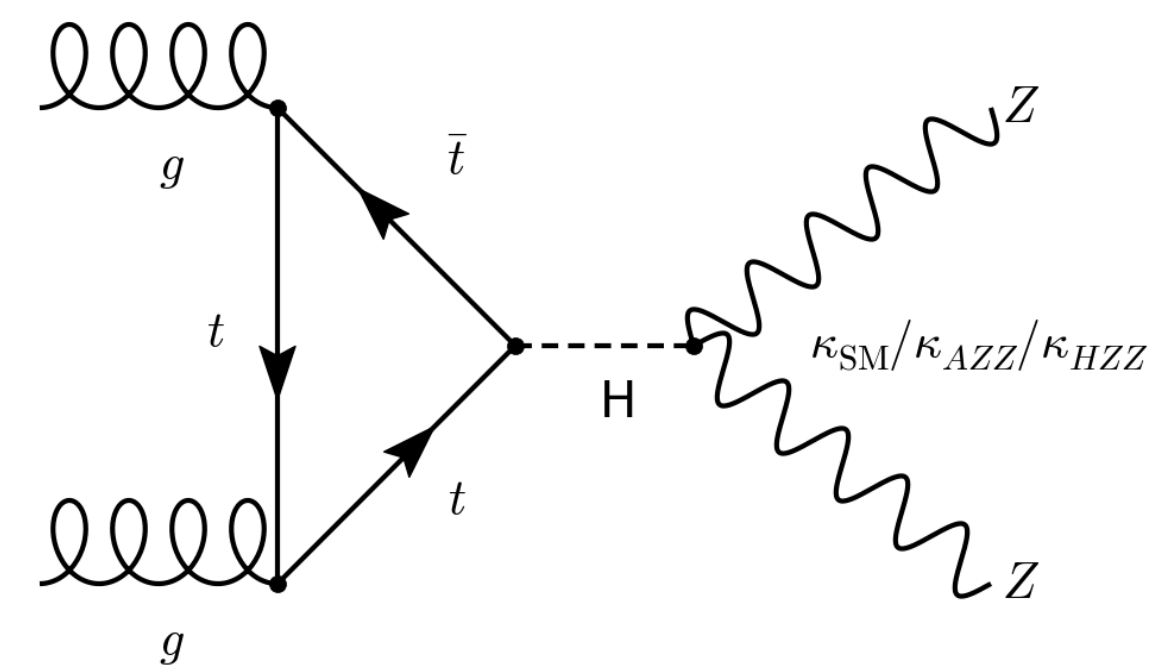
- Signal samples
 - The signal is modelled using the POWHEG Monte Carlo event generator at NLO.
 - QCD corrections up to NLO and QCD soft-gluon re-summations up to next-to-next-to-leading logarithm are applied.
 - The Higgs boson decay branching ratio in this final state is predicted by Prophecy4f, with includes NLO QCD+EW corrections, the interferences effects between identical final states and the leading two-loop heavy Higgs boson corrections to the four-fermion width.
- Main Background
 - continuum $((Z^{(*)}/\gamma^{(*)})(Z^{(*)}/\gamma^{(*)})$ production
 - Sherpa 2.2.1 generator
 - Matrix elements contain all diagrams with four electroweak vertices. They are calculated for up to 1 parton at NLO and up to 3 partons at LO



Bonus:

CP study in the decay vertex

$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{HVV}$ AND κ_{AVV} FIT



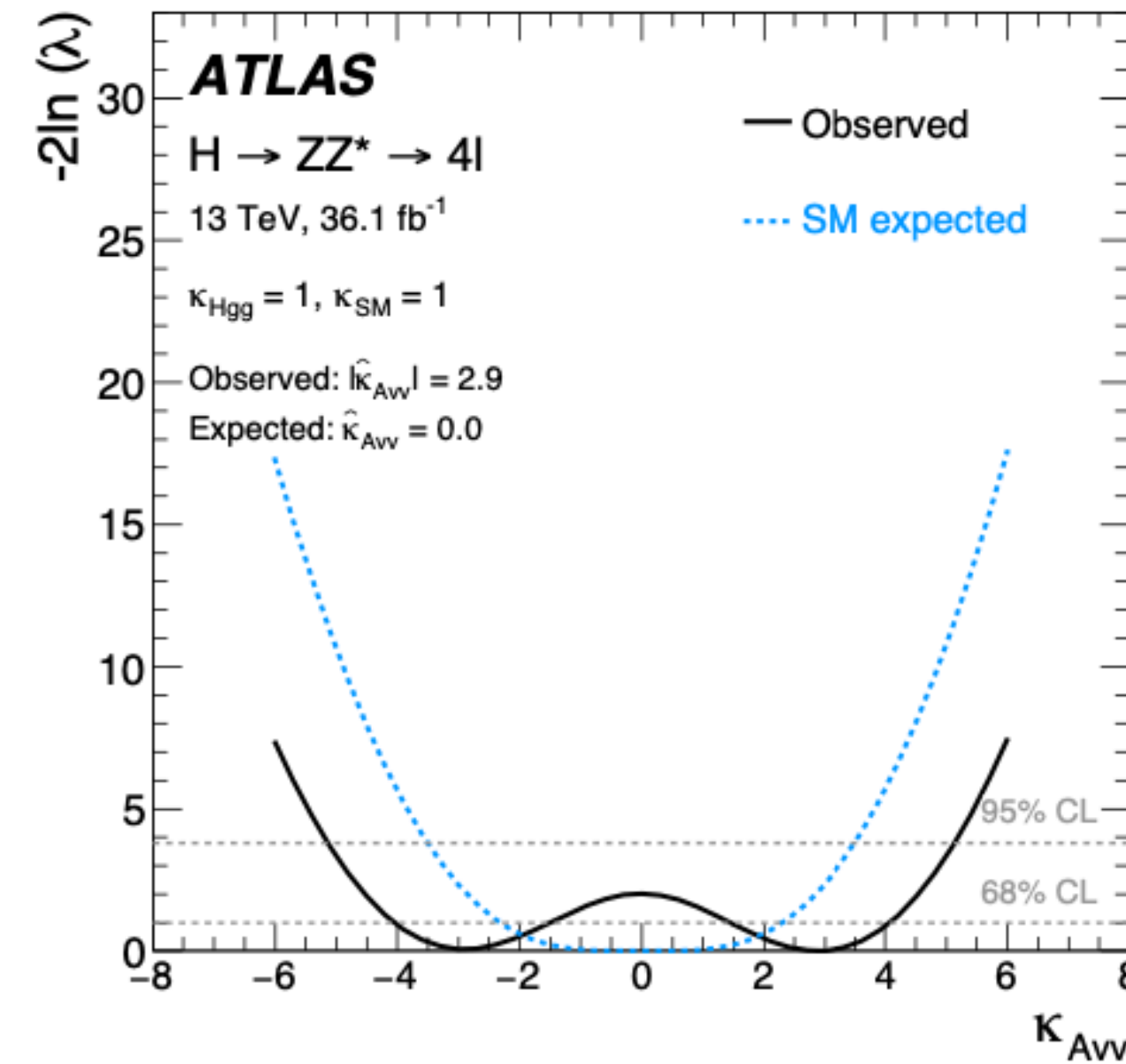
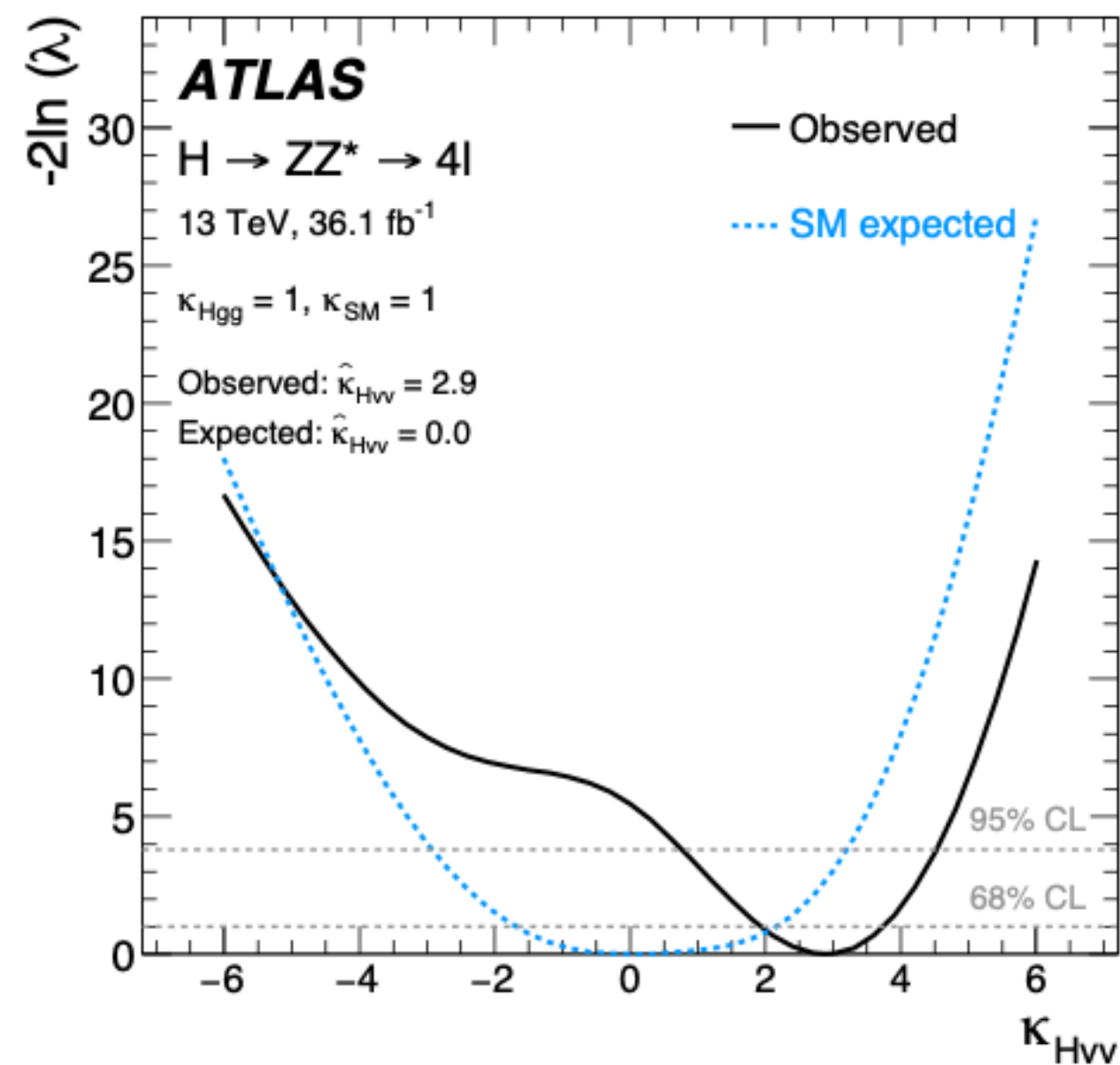
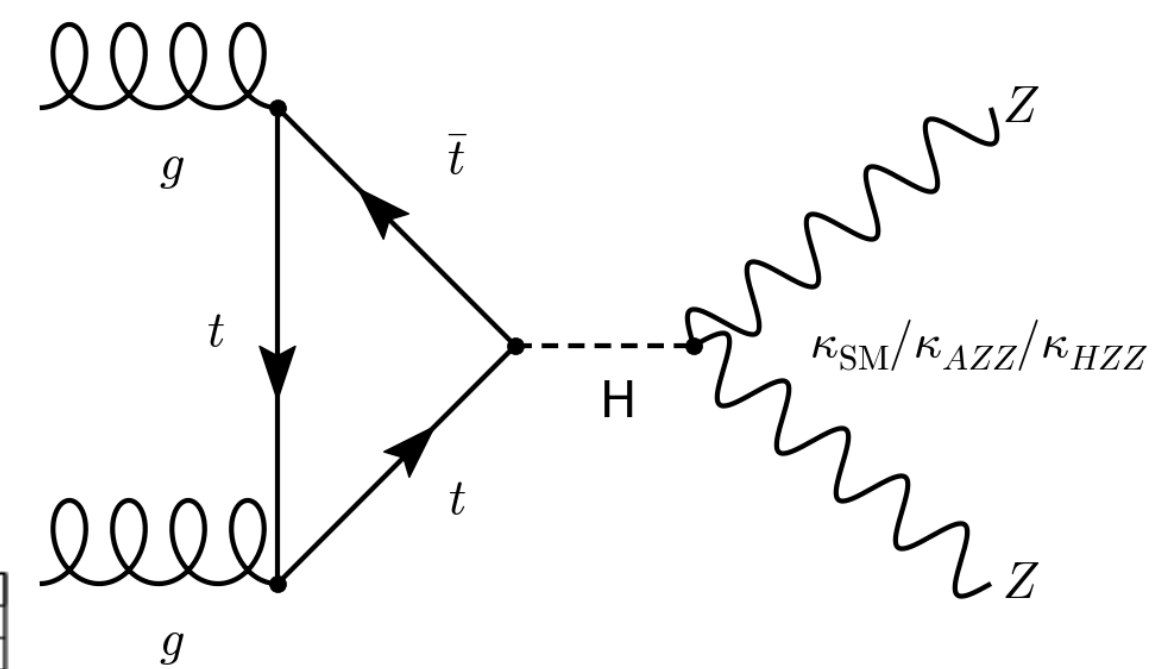
$$\mathcal{L}_0^V =$$

$$-\frac{1}{4} \frac{1}{\Lambda} [\kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + \tan \alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu}]$$

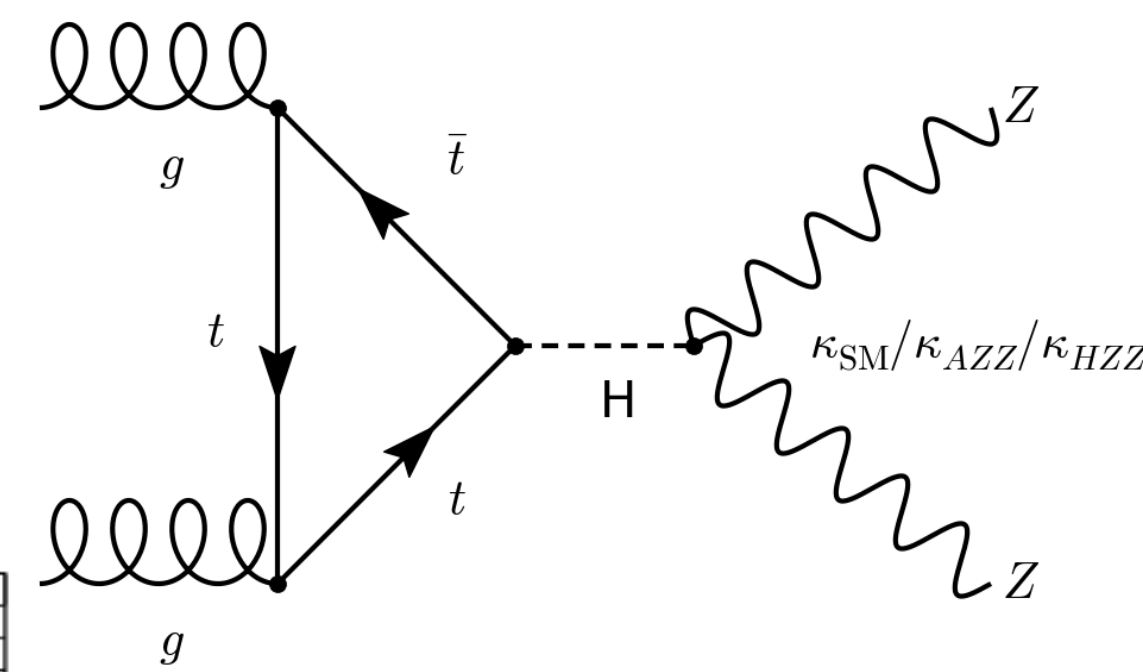
$$-\frac{1}{2} \frac{1}{\Lambda} [\kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + \tan \alpha \kappa_{AWW} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu}] \} \mathcal{X}_0$$

- Probing BSM couplings in the HVV vertex
 - The total width of the Higgs boson is modified to take into account the changes in the BR due to the BSM contributions

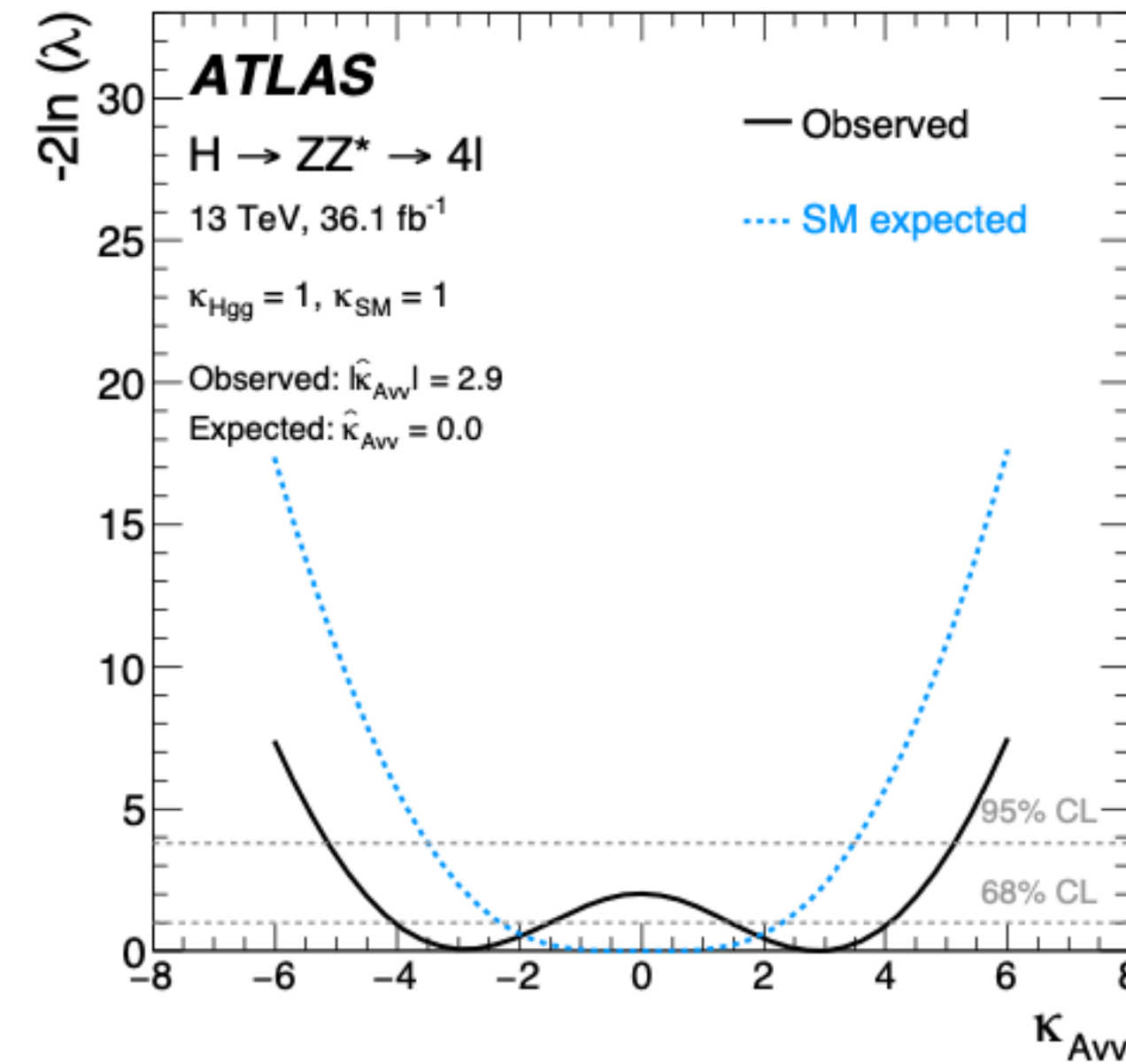
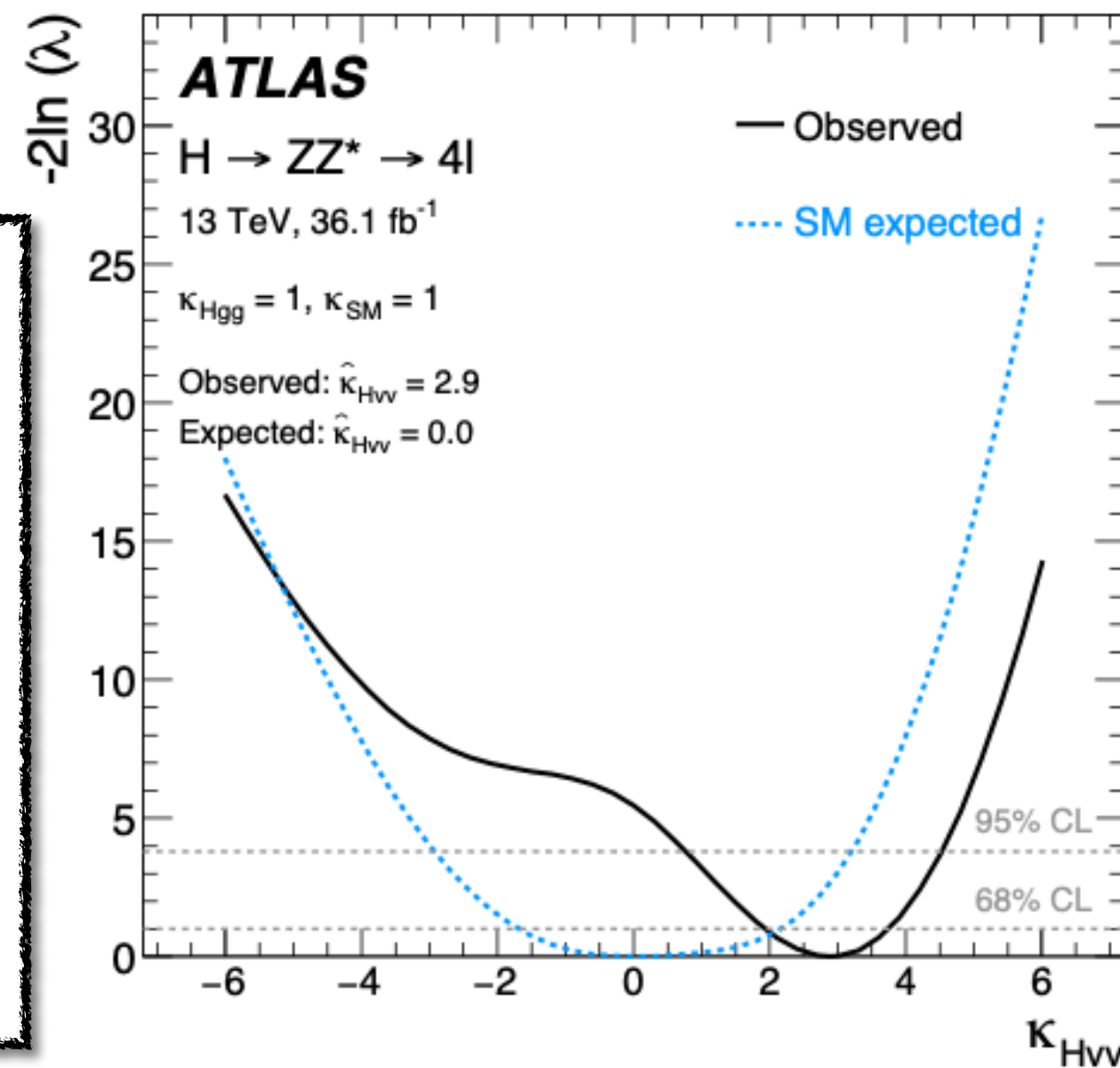
$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{HVV}$ AND κ_{AVV} FIT



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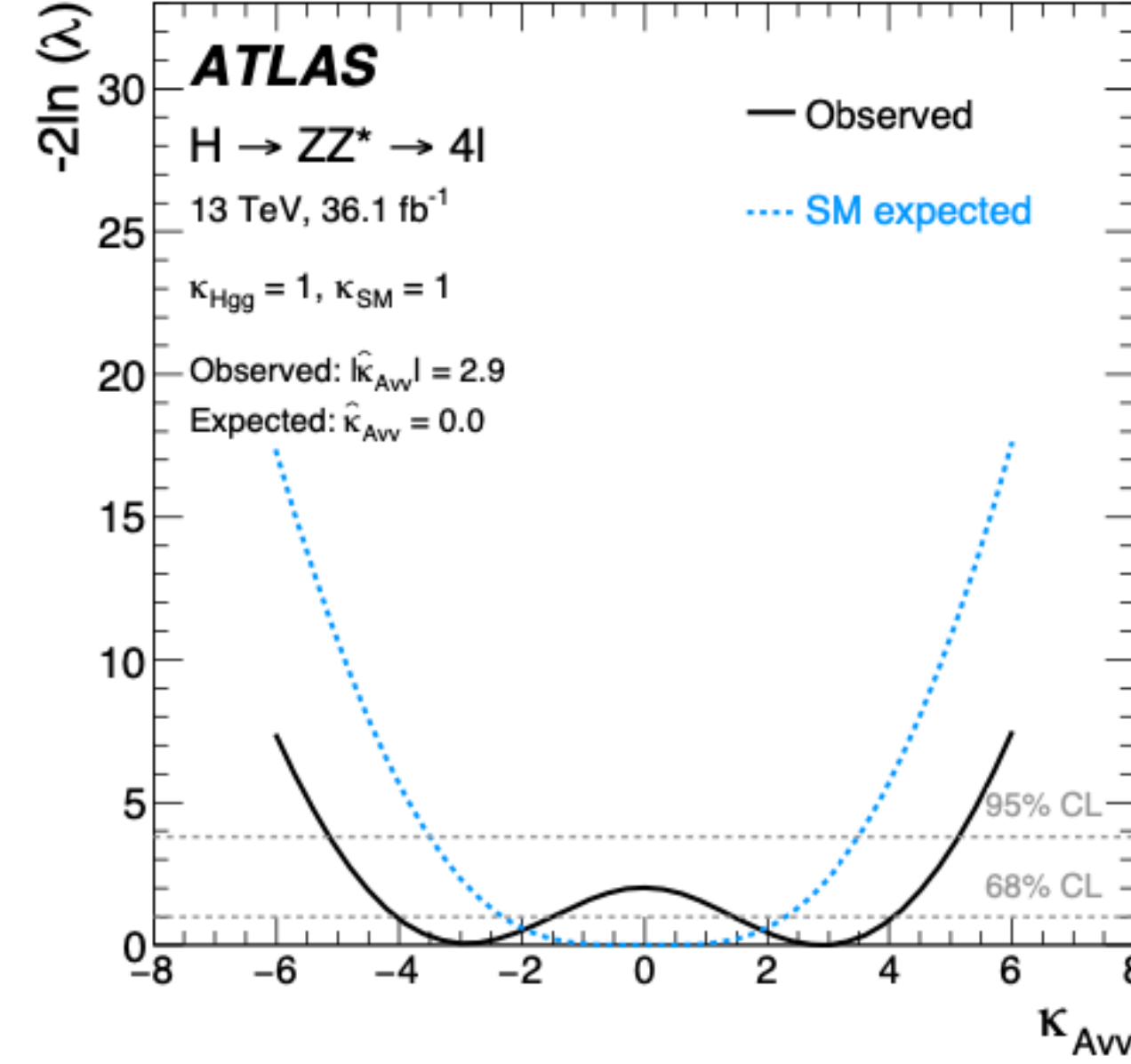
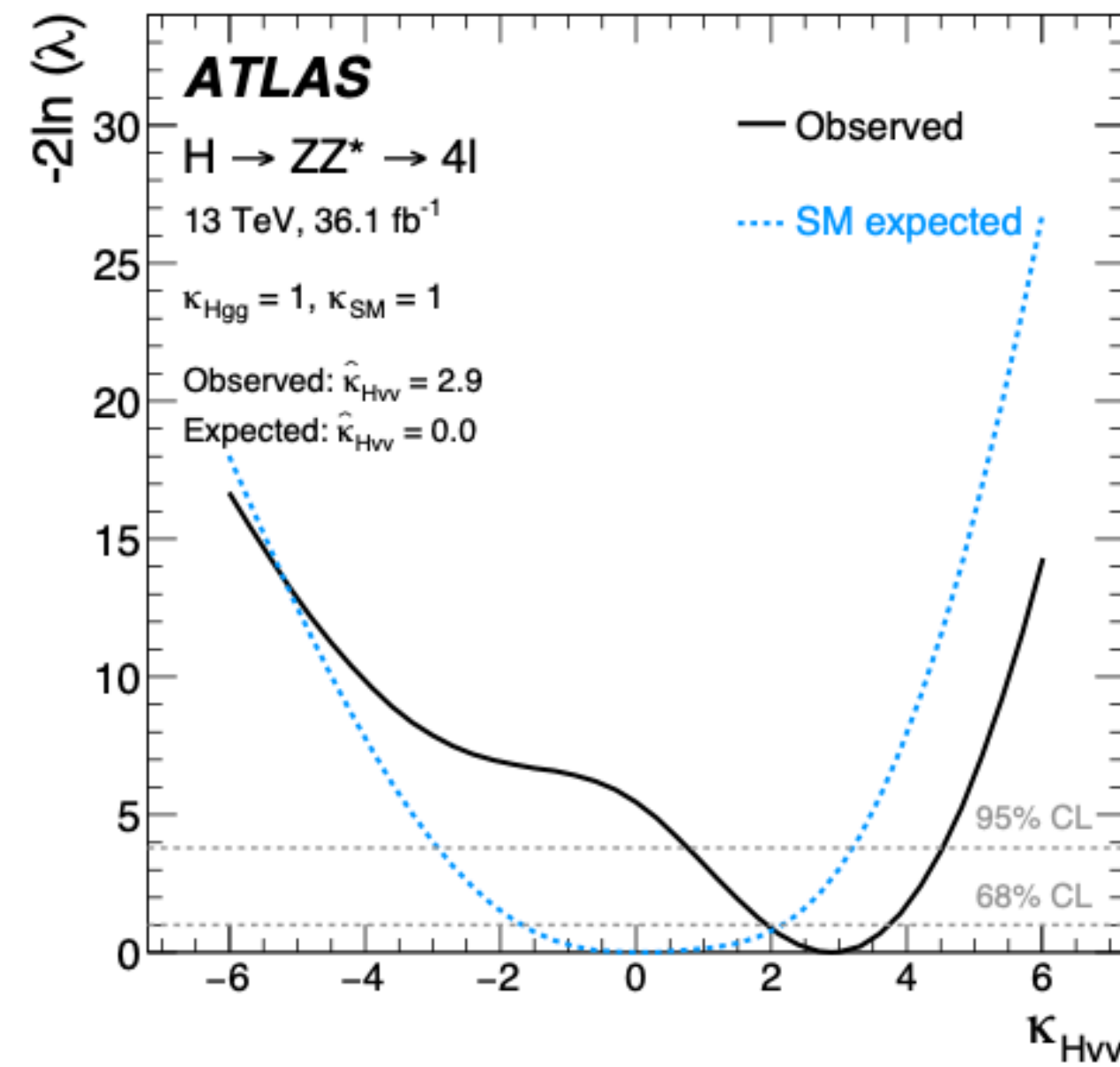
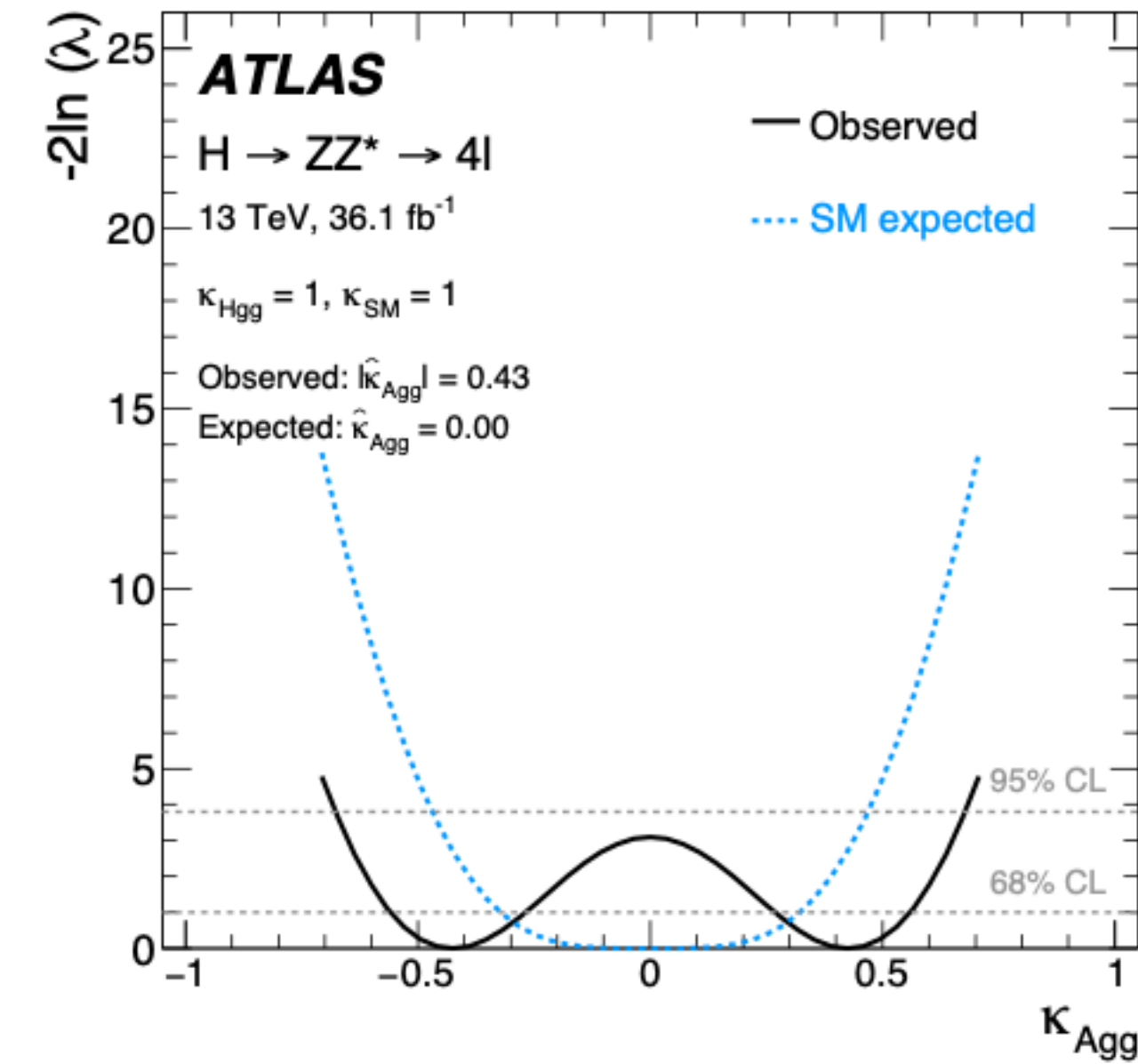
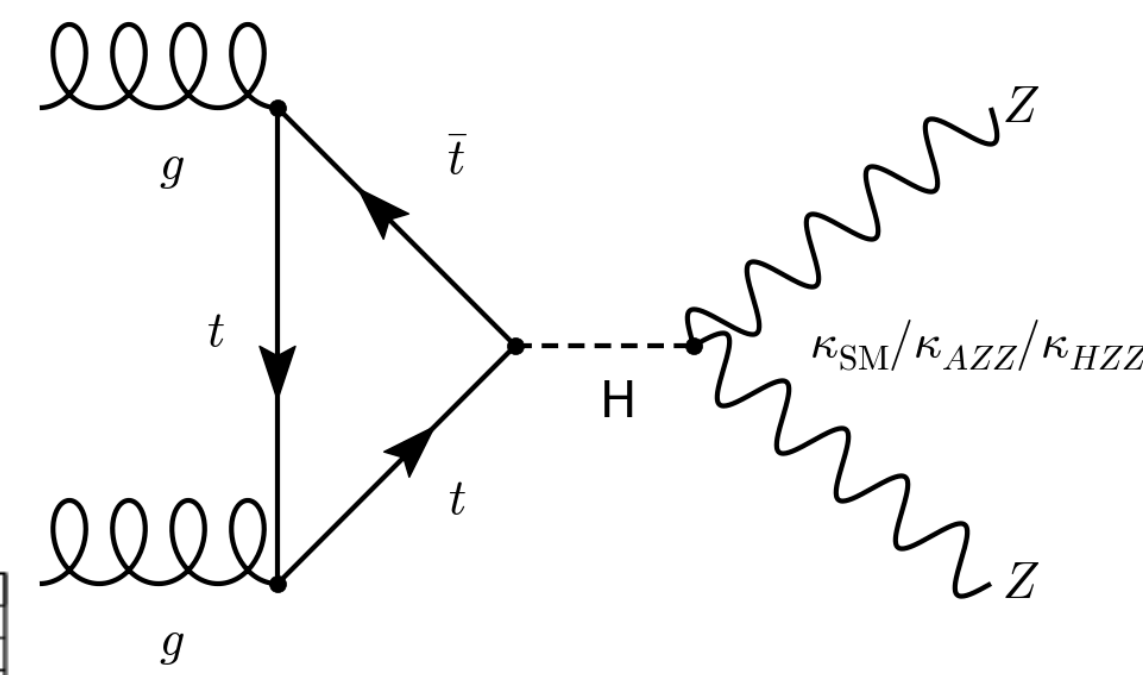


Asymmetry is expected due to the interference between SM and BSM contribution



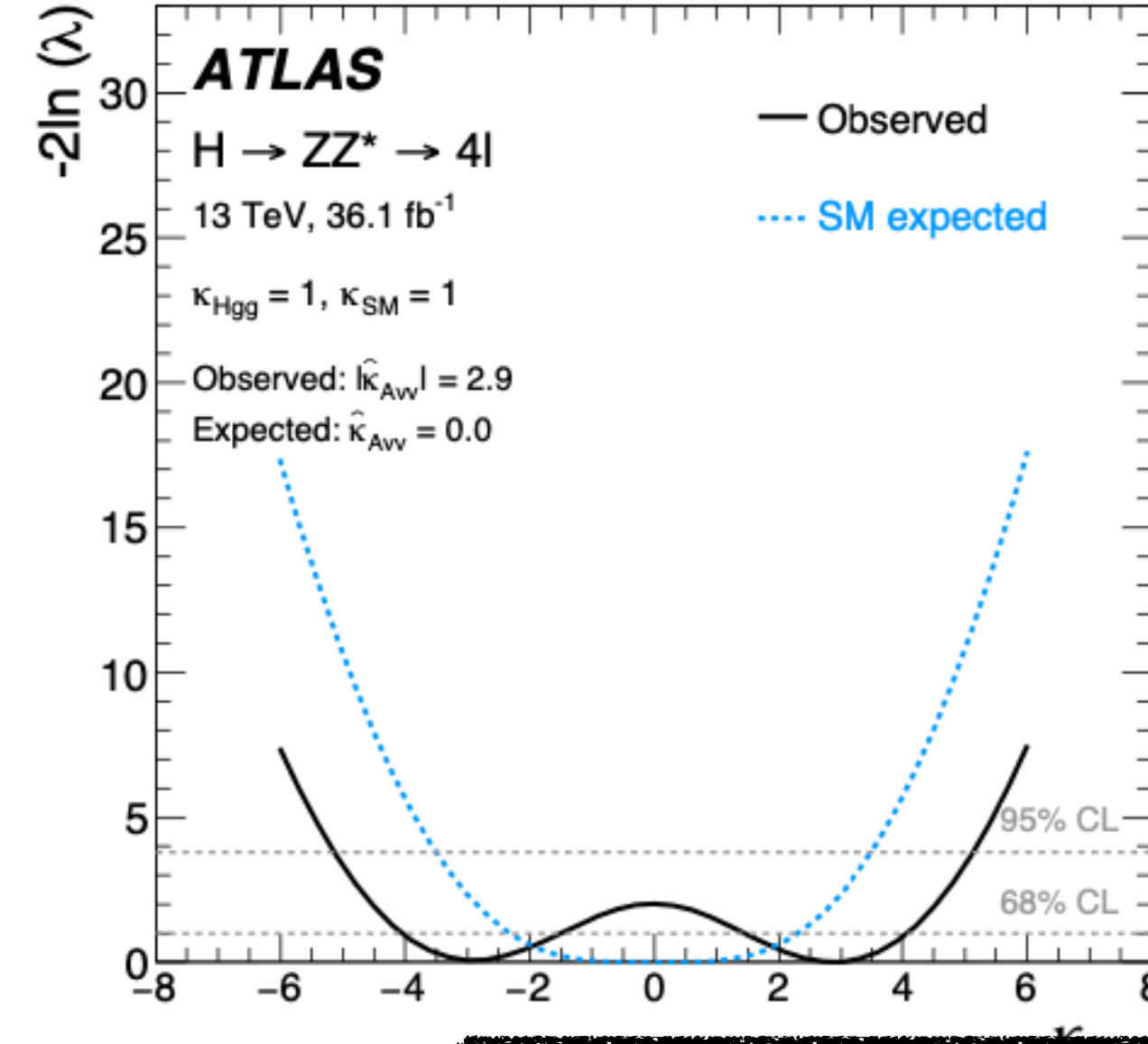
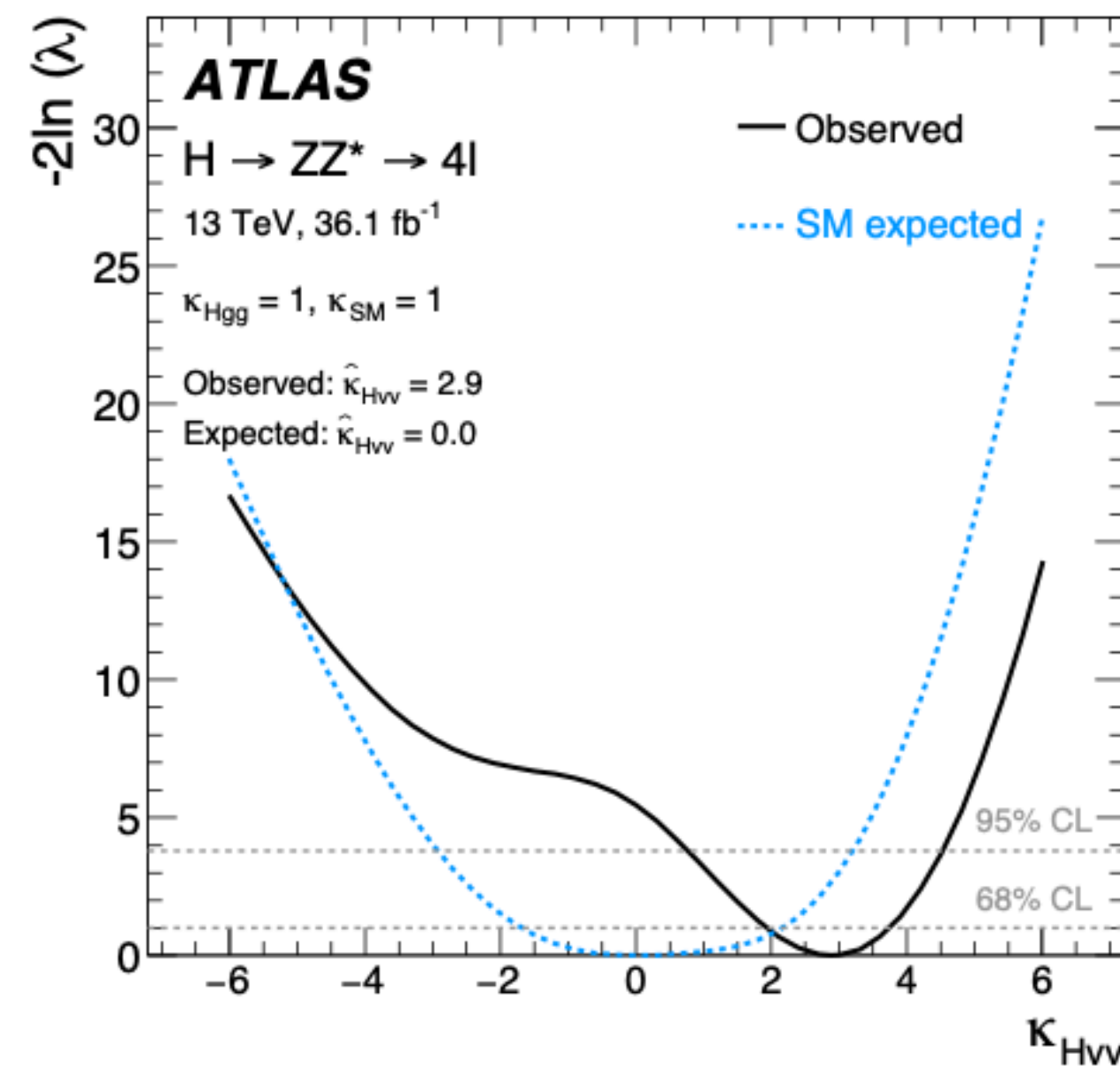
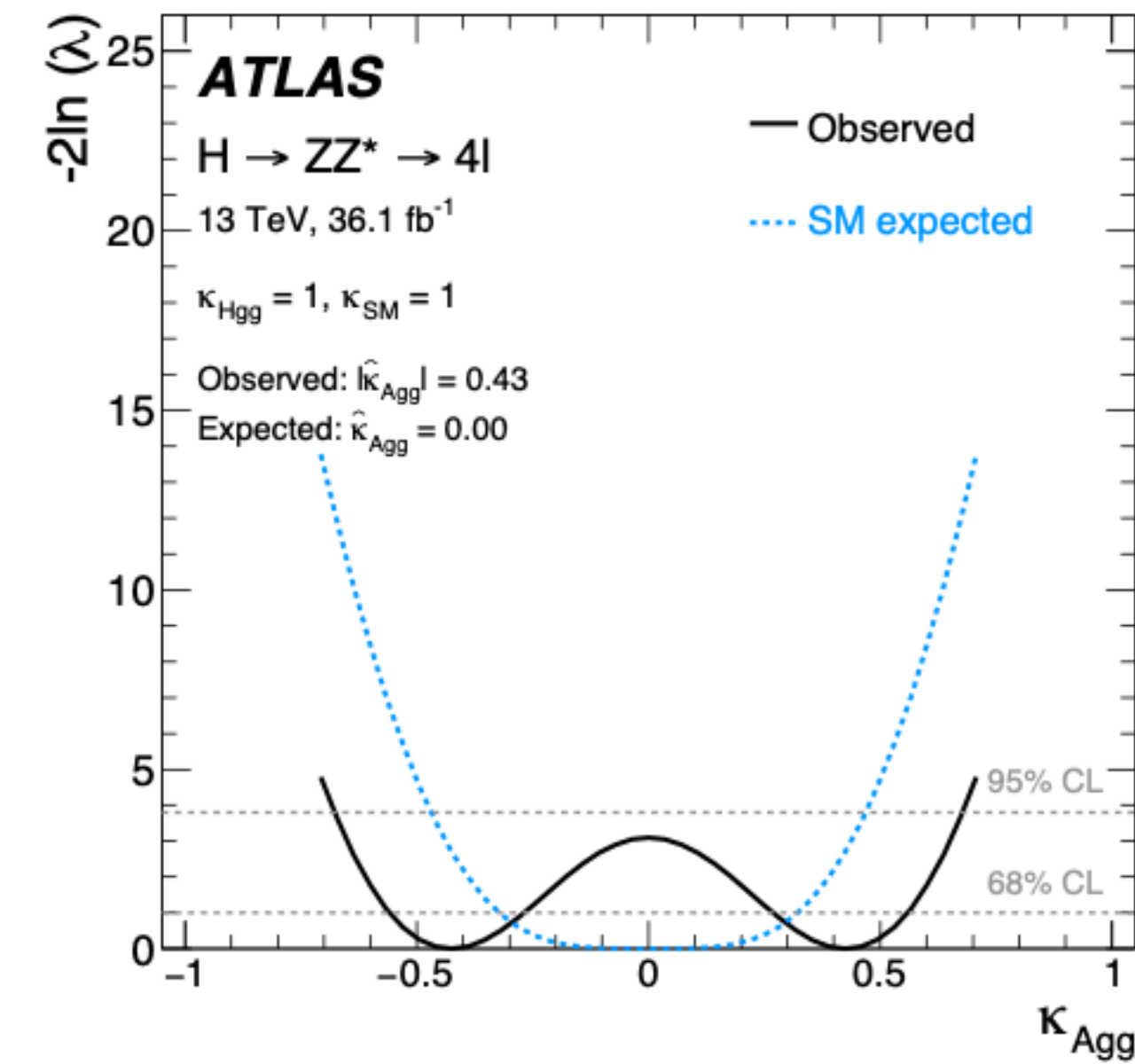
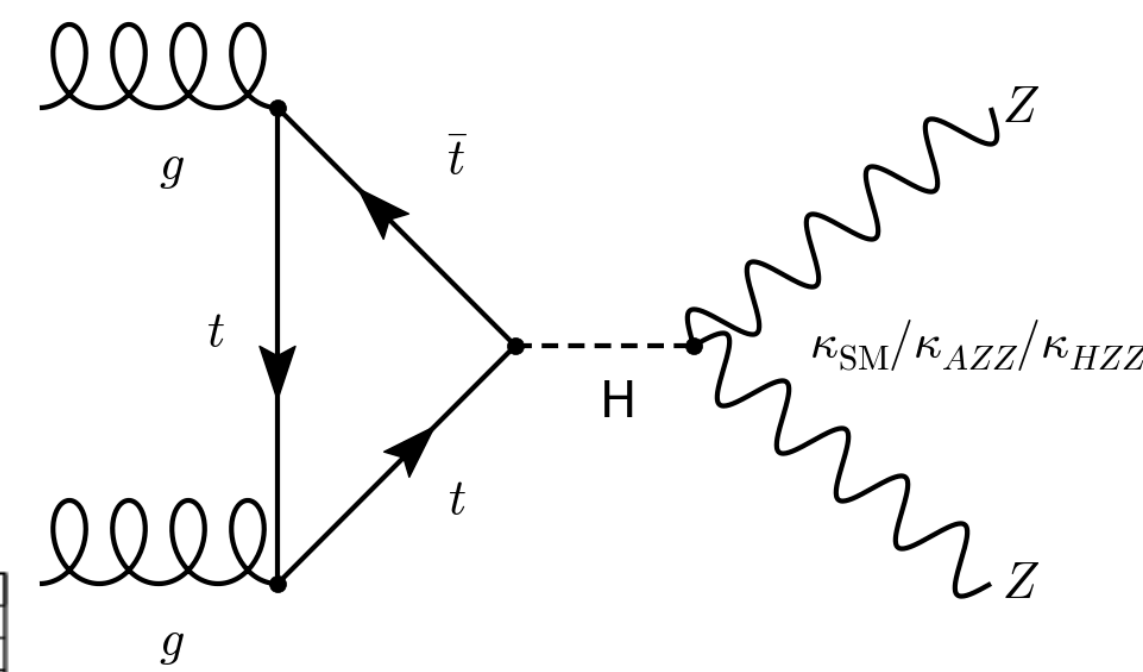
It is symmetric because the cross section scales at the same rate for negative and positive values of κ_{AVV}

H \rightarrow ZZ* \rightarrow 4l, κ_{Agg} , κ_{HVV} AND κ_{AVV} FIT



BSM coupling	Fit configuration	Expected conf. inter.	Observed conf. inter.	Best-fit $\hat{\kappa}_{BSM}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
κ_{Agg}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-0.47, 0.47]$	$[-0.68, 0.68]$	± 0.43	-	1.8σ
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-2.9, 3.2]$	$[0.8, 4.5]$	2.9	-	2.3σ
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	$[-3.1, 4.0]$	$[-0.6, 4.2]$	2.2	1.2	1.7σ
κ_{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-3.5, 3.5]$	$[-5.2, 5.2]$	± 2.9	-	1.4σ
κ_{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	$[-4.0, 4.0]$	$[-4.4, 4.4]$	± 1.5	1.2	0.5σ

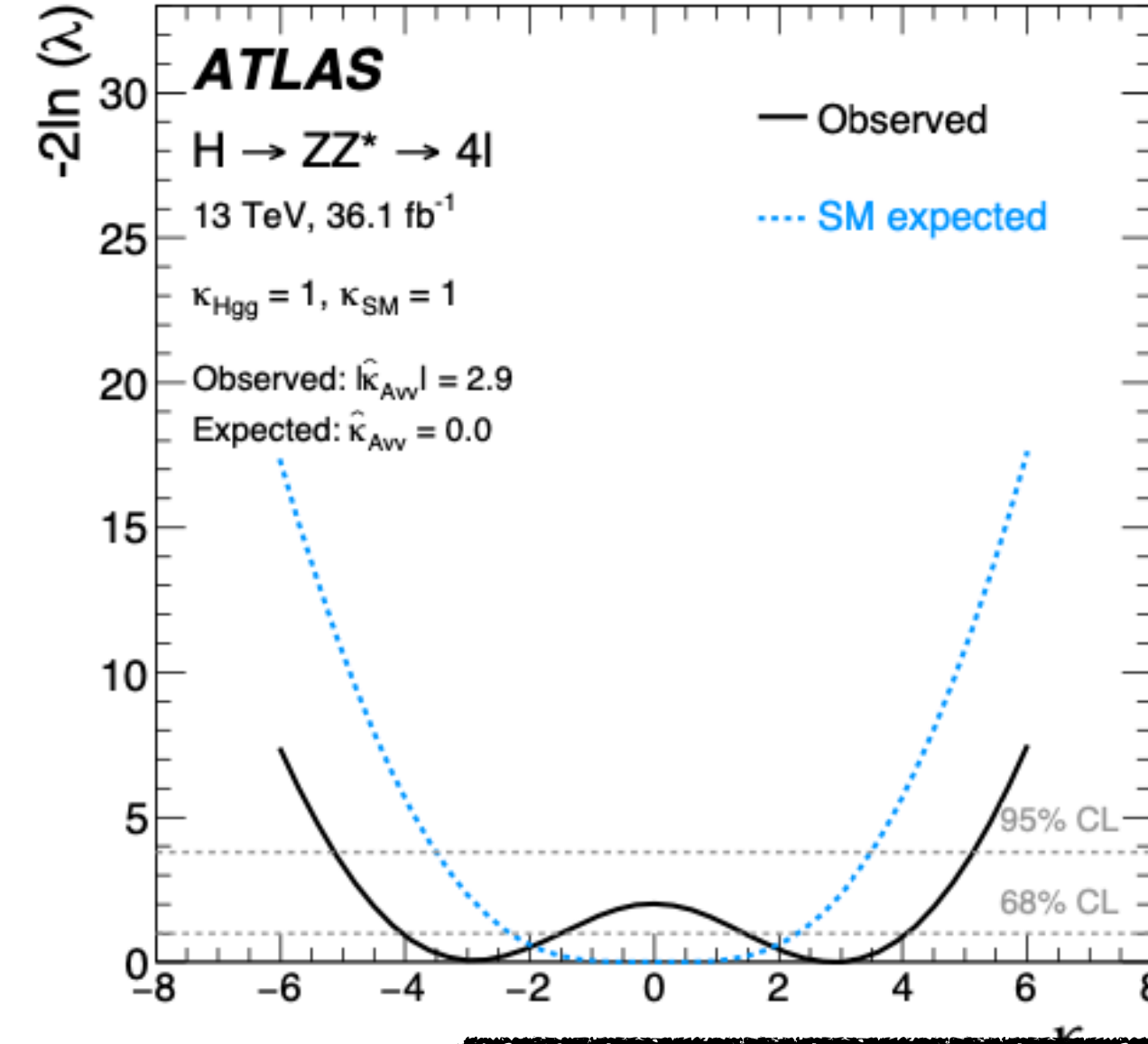
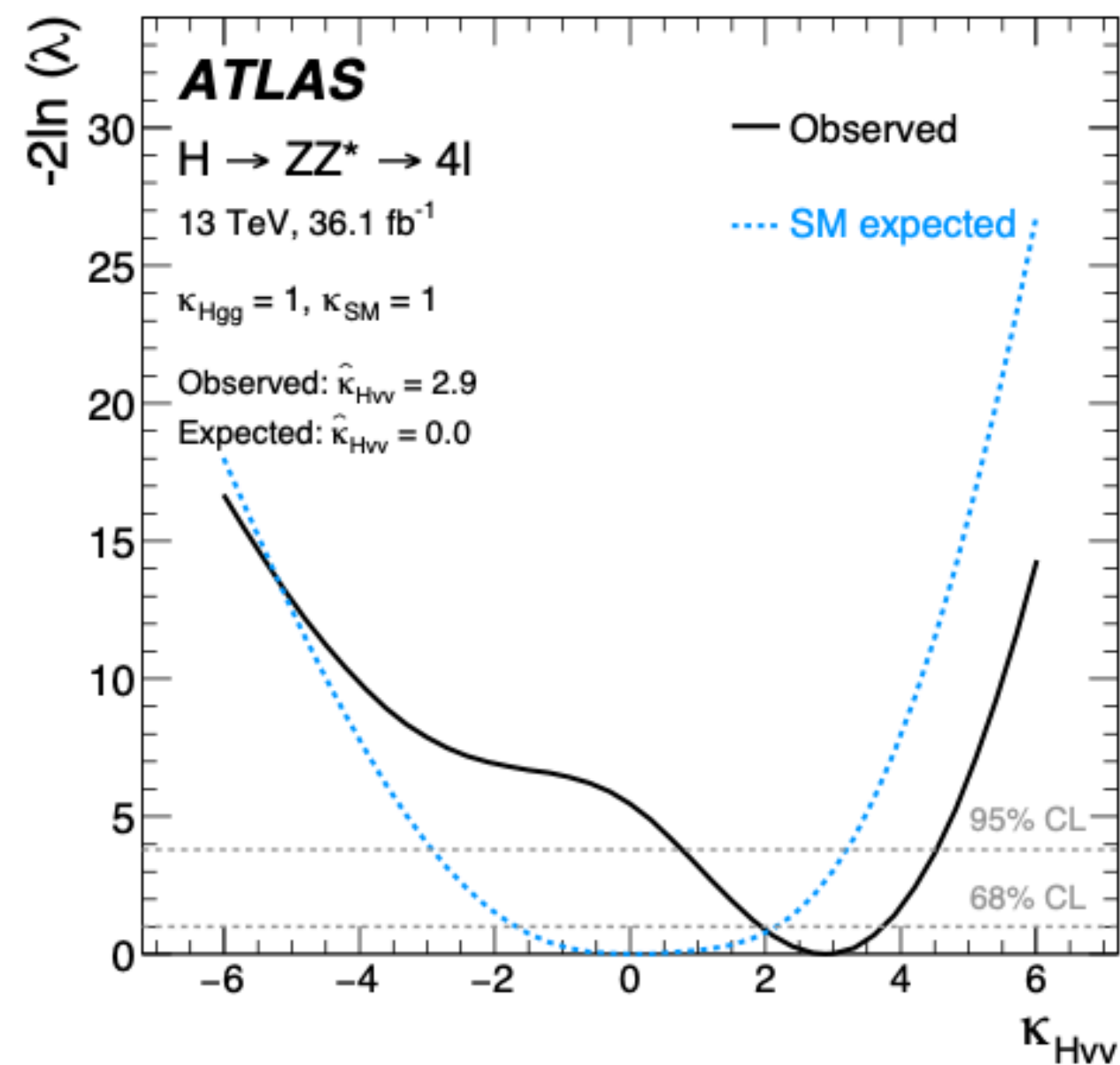
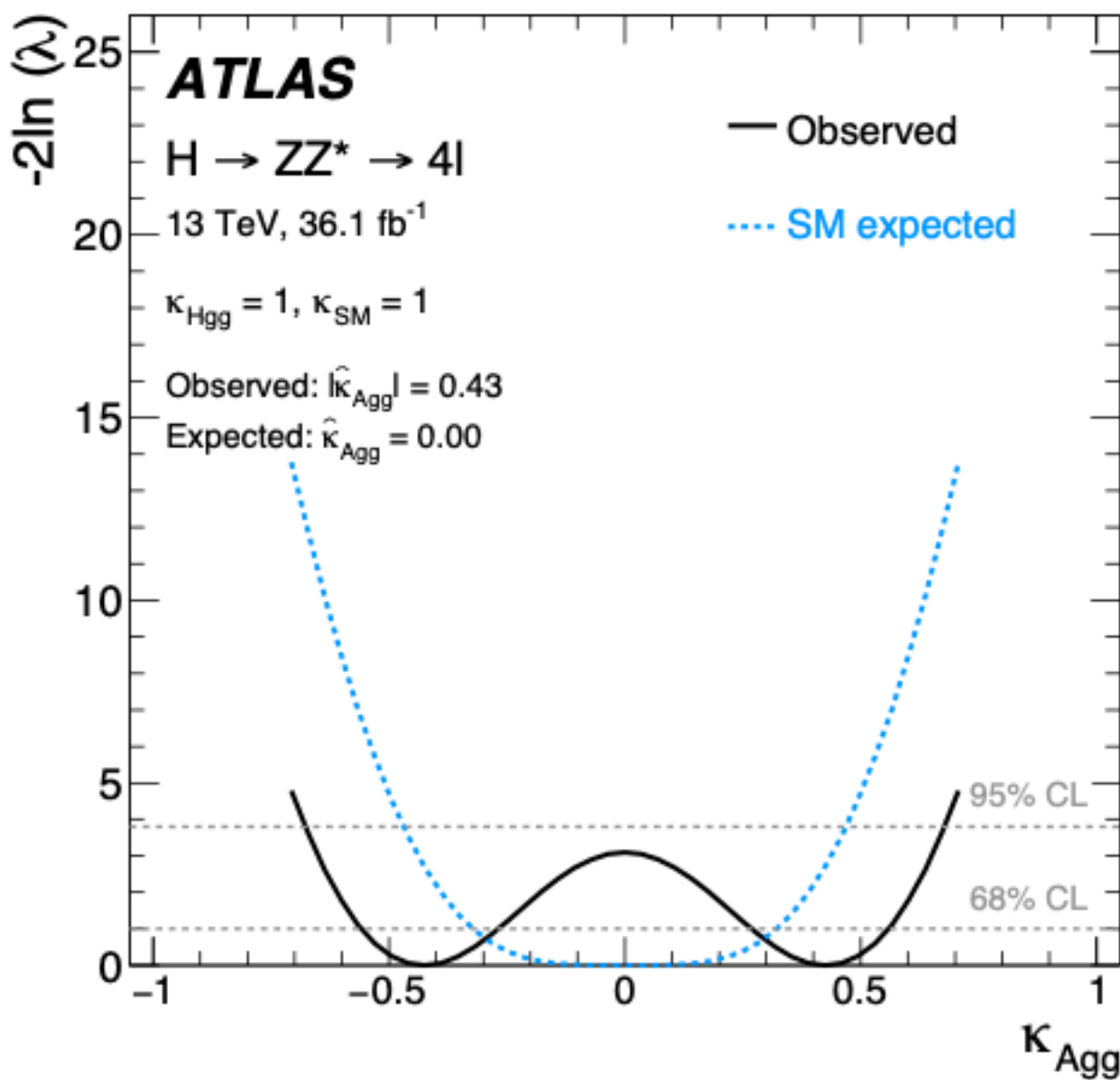
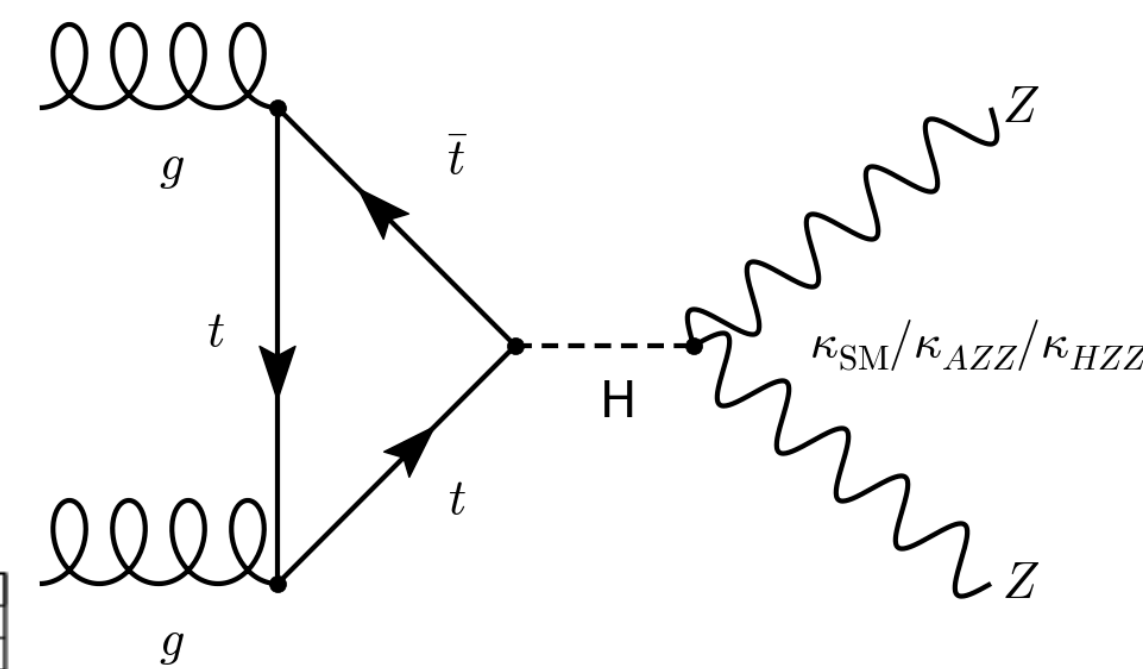
H \rightarrow ZZ* \rightarrow 4l, κ_{Agg} , κ_{HVV} AND κ_{AVV} FIT



BSM coupling	Fit configuration	Expected conf. inter.	Observed conf. inter.	Best-fit $\hat{\kappa}_{BSM}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
κ_{Agg}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-0.47, 0.47]$	$[-0.68, 0.68]$	± 0.43	-	1.8σ
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-2.9, 3.2]$	$[0.8, 4.5]$	2.9	-	2.3σ
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	$[-3.1, 4.0]$	$[-6, 4.2]$	2.2	1.2	1.7σ
κ_{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-3.5, 3.5]$	$[-5.2, 5.2]$	± 2.9	-	1.4σ
κ_{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	$[-4.0, 4.0]$	$[-4.4, 4.4]$	± 1.5	1.2	0.5σ

If κ_{SM} is left free in the fit, the expected limits on the BSM couplings decrease by up to 10%

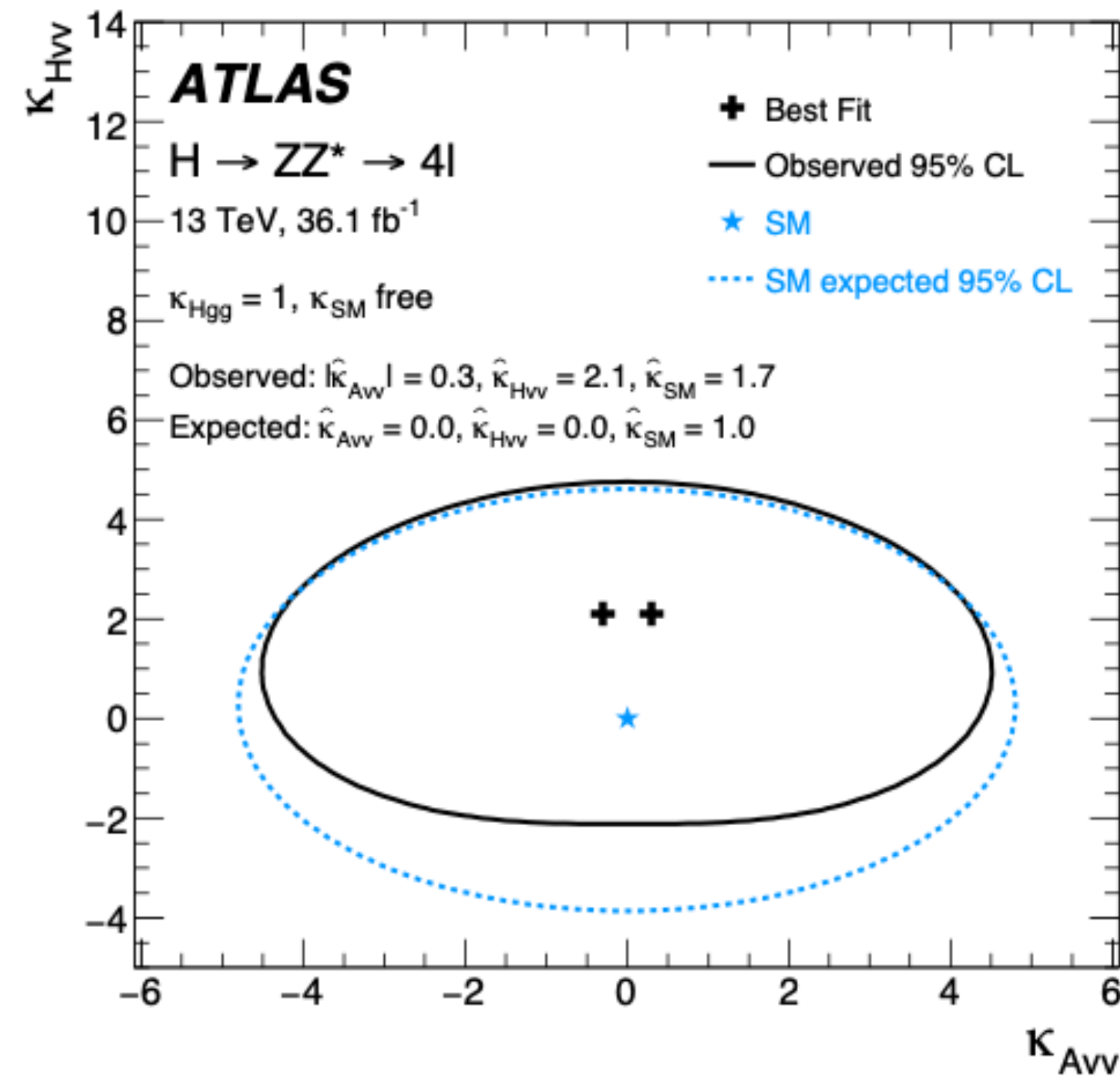
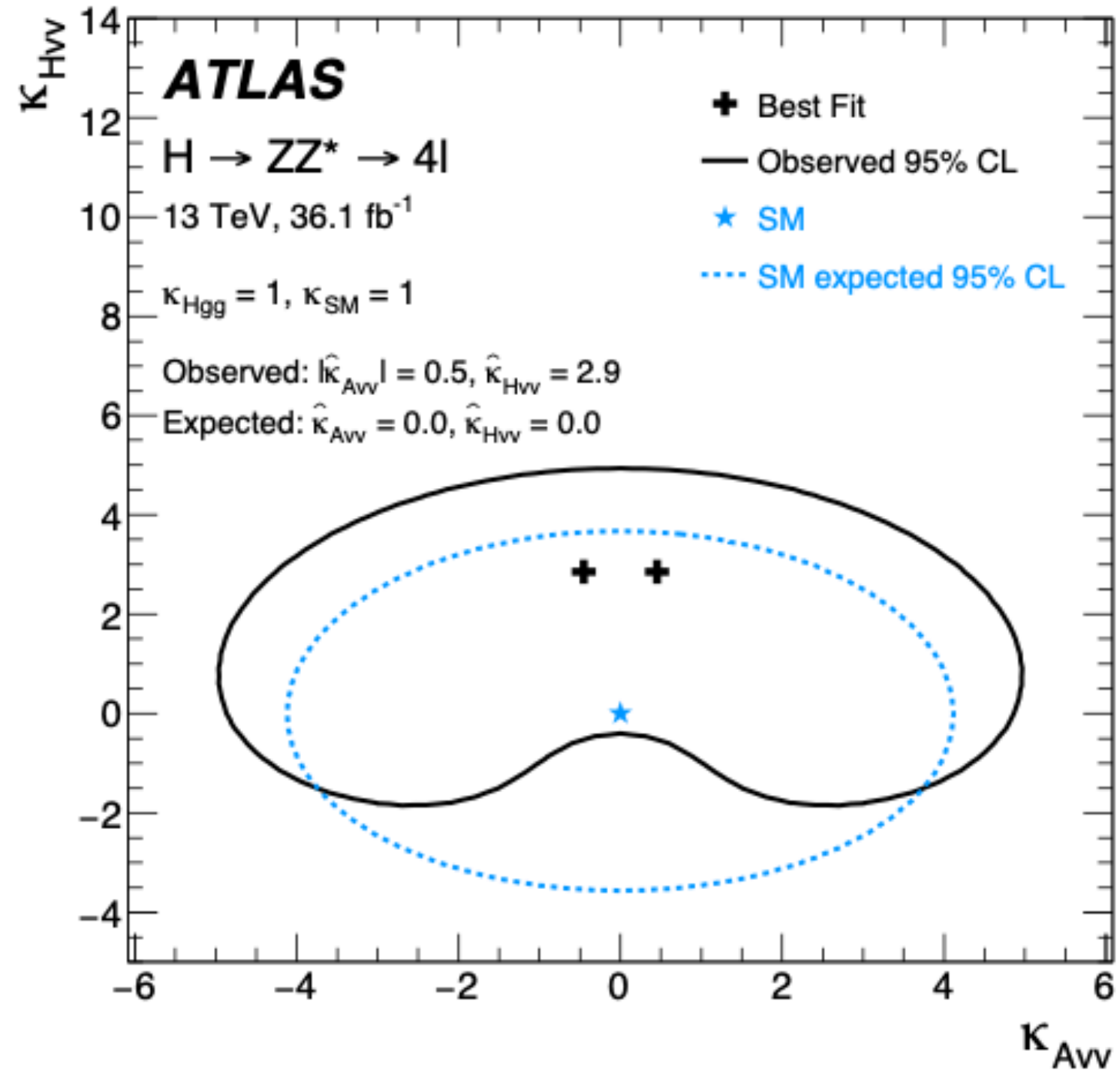
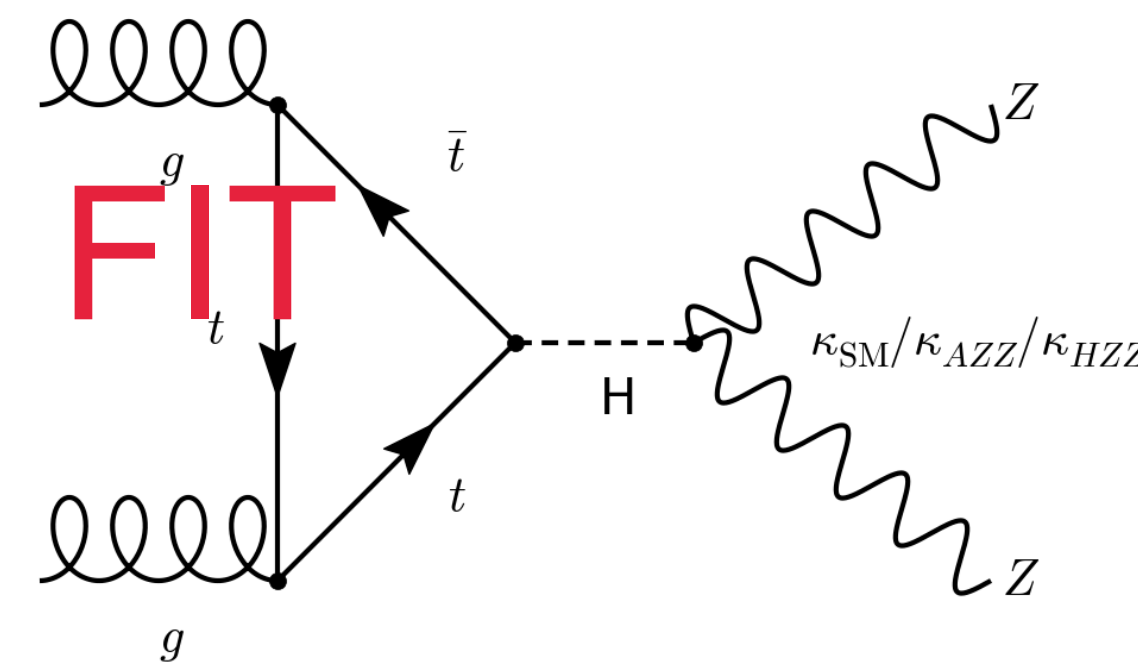
$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{Agg}, \kappa_{HVV}$ AND κ_{AVV} FIT



BSM coupling	Fit configuration	Expected conf. inter.	Observed conf. inter.	Best-fit $\hat{\kappa}_{BSM}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
κ_{Agg}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-0.47, 0.47]$	$[-0.68, 0.68]$	± 0.43	-	1.8σ
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-2.9, 3.2]$	$[0.8, 4.5]$	2.9	-	2.3σ
κ_{HVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	$[-3.1, 4.0]$	$[-0.6, 4.2]$	2.2	1.2	1.7σ
κ_{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} = 1)$	$[-3.5, 3.5]$	$[-5.2, 5.2]$	± 2.9	-	1.4σ
κ_{AVV}	$(\kappa_{Hgg} = 1, \kappa_{SM} \text{ free})$	$[-4.0, 4.0]$	$[-4.4, 4.4]$	± 1.5	1.2	0.5σ

The deviation of the BSM couplings from the SM decreases because of the increase of the value of $\hat{\kappa}_{SM}$

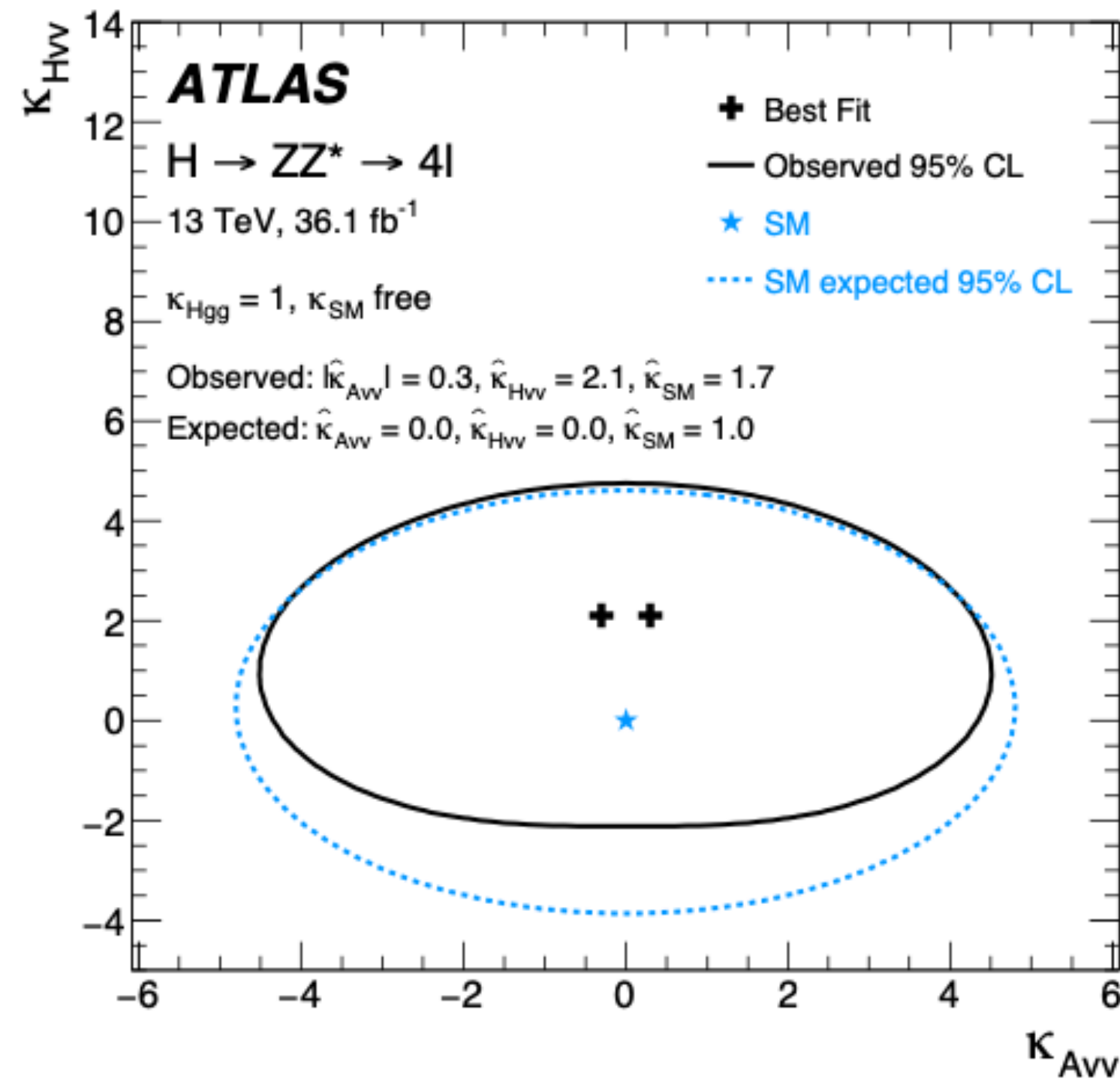
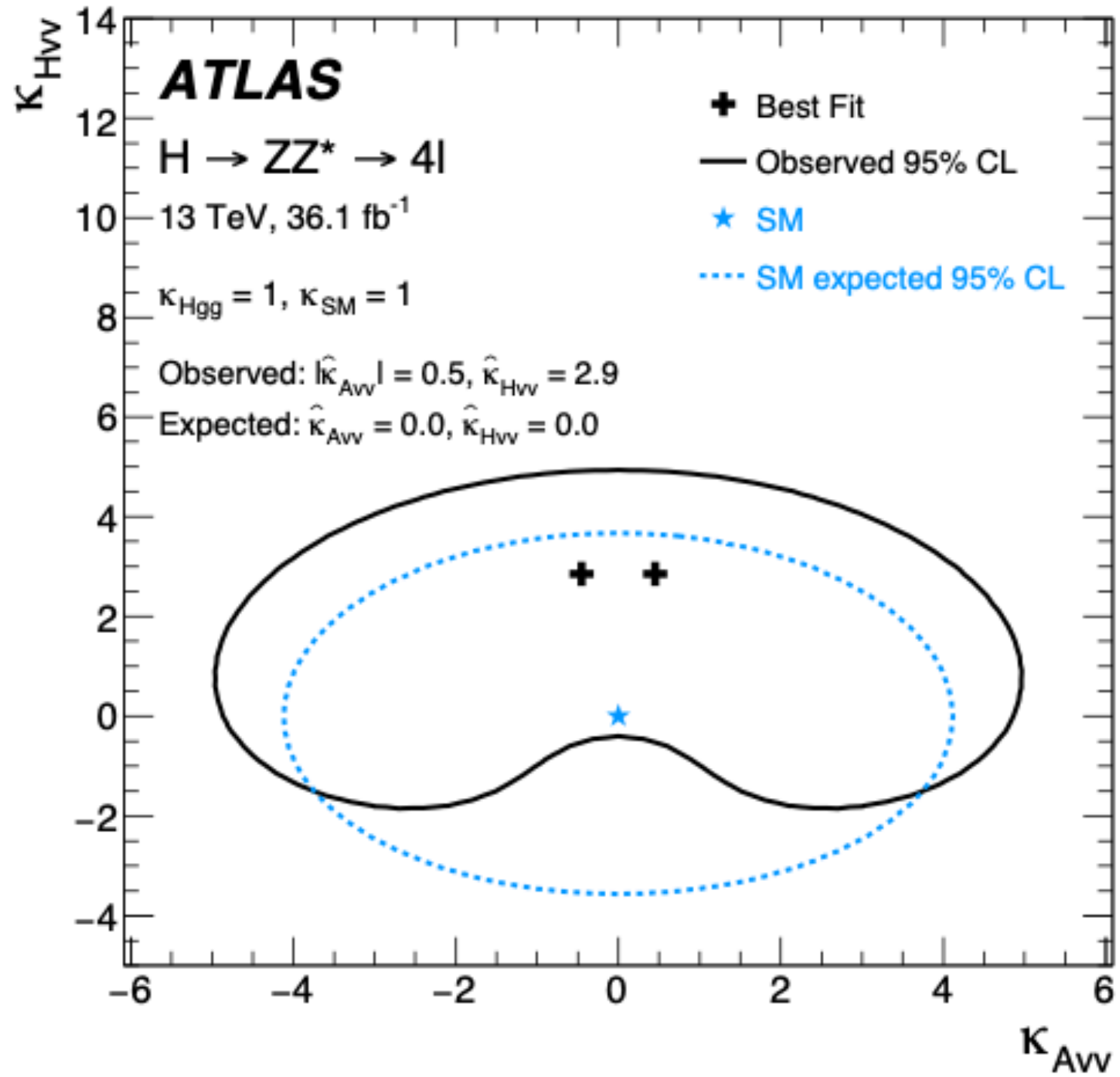
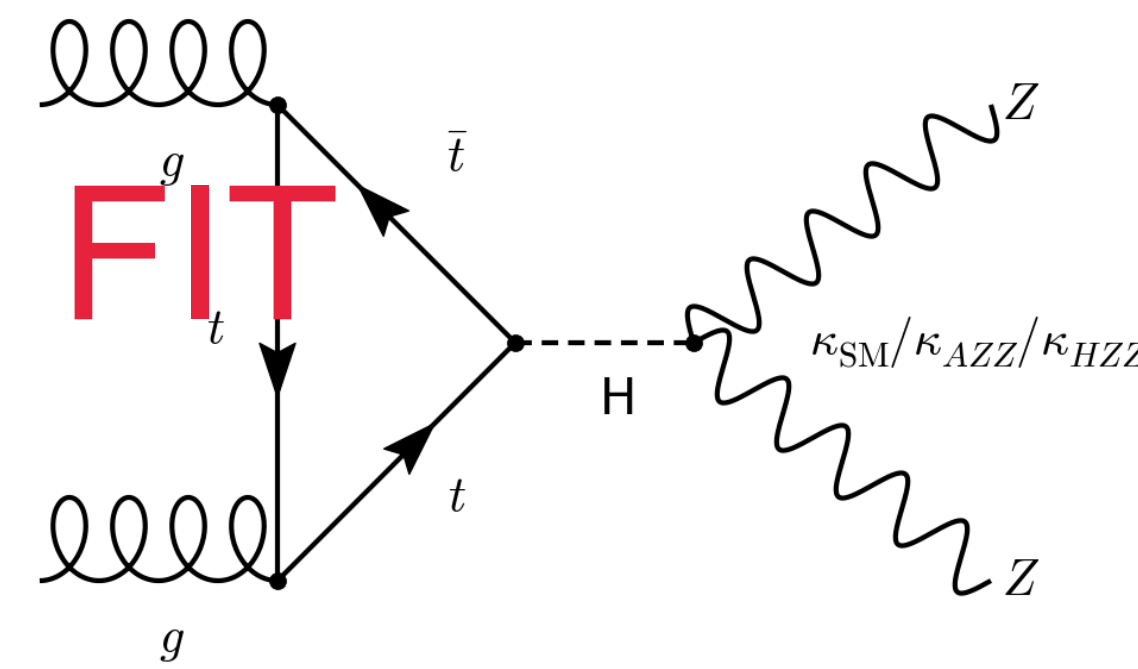
$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{HVV}, \kappa_{AVV}$ SIMULTANEOUS FIT



Asymmetric because of κ_{HVV}

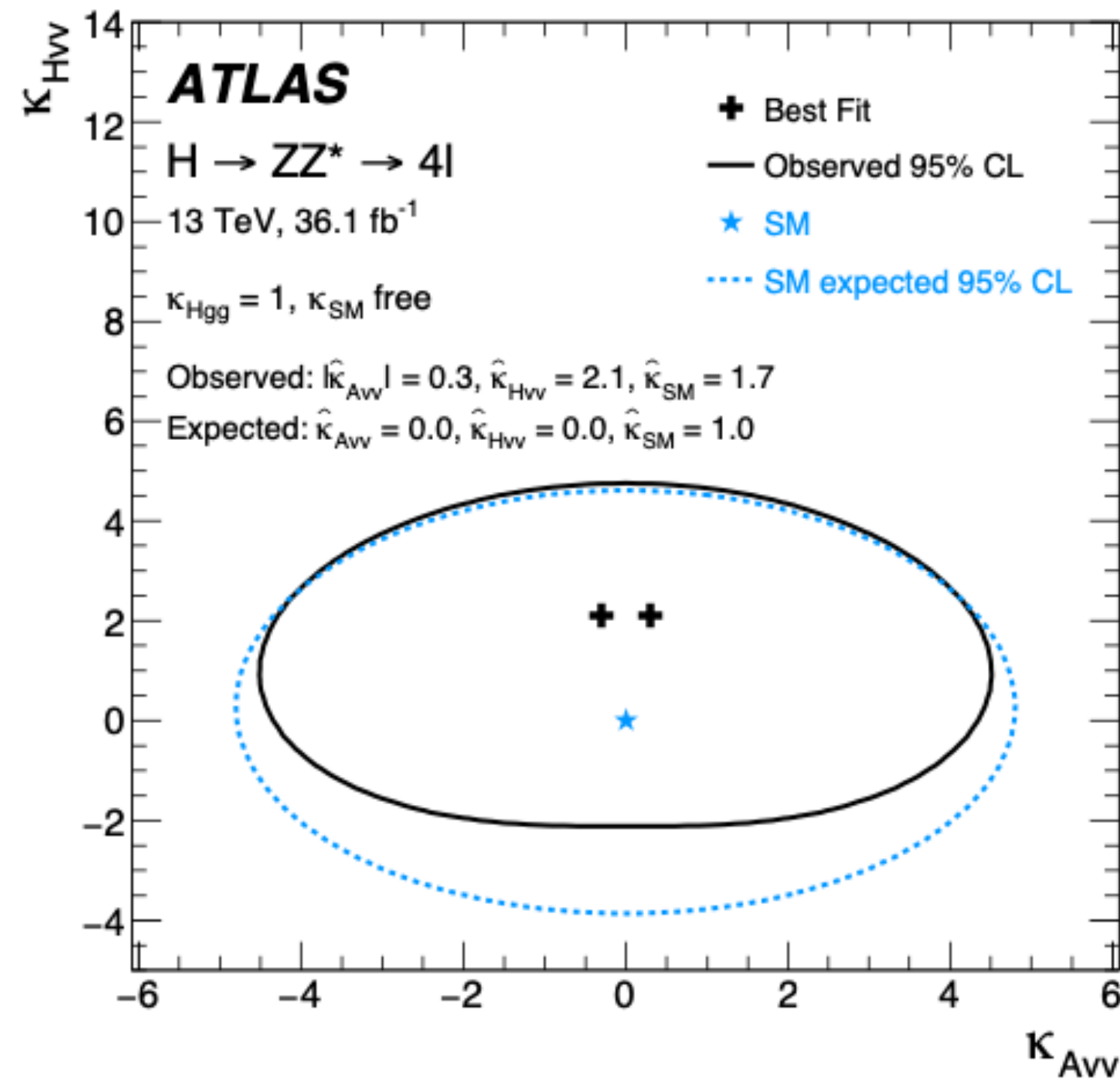
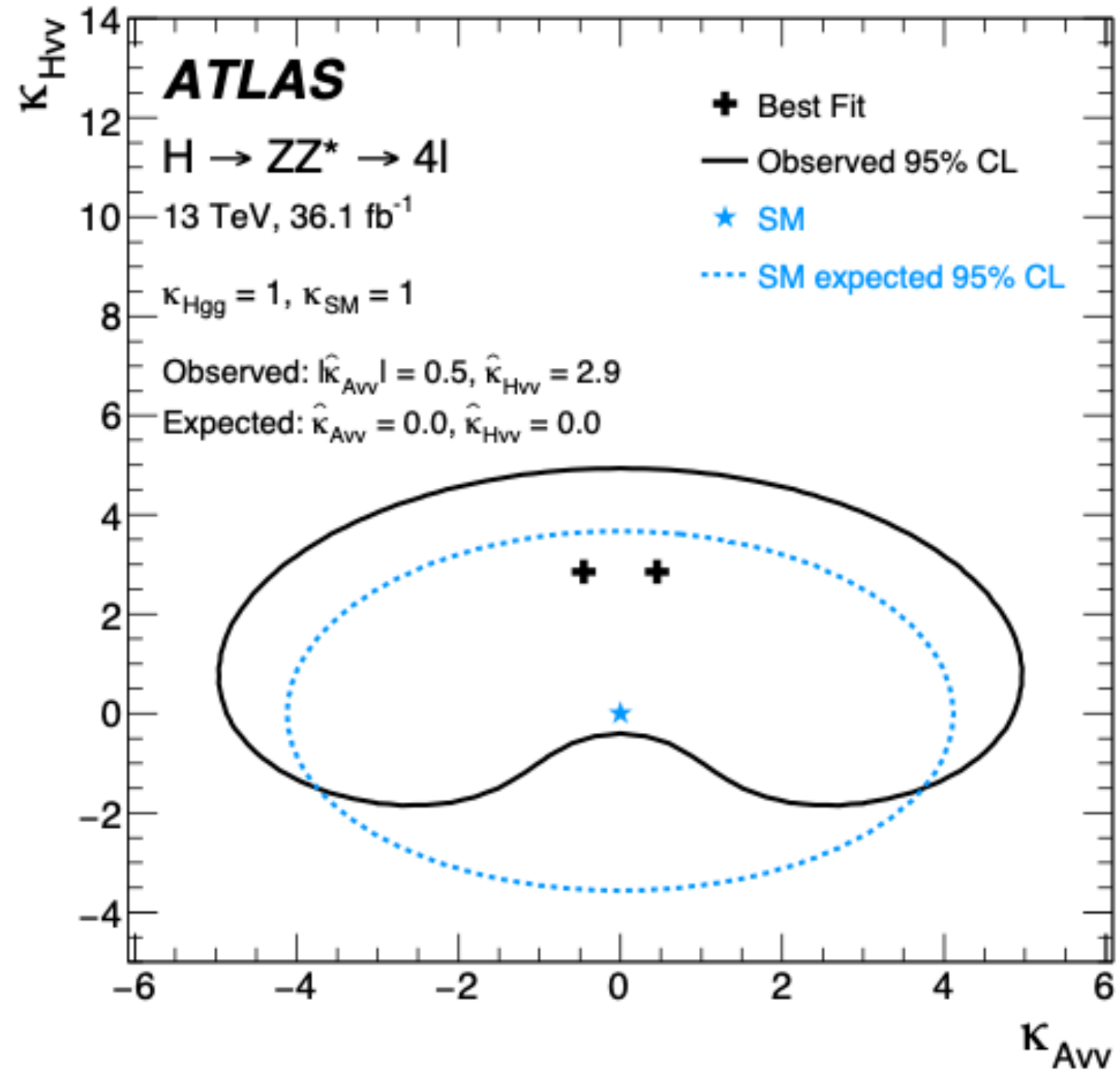
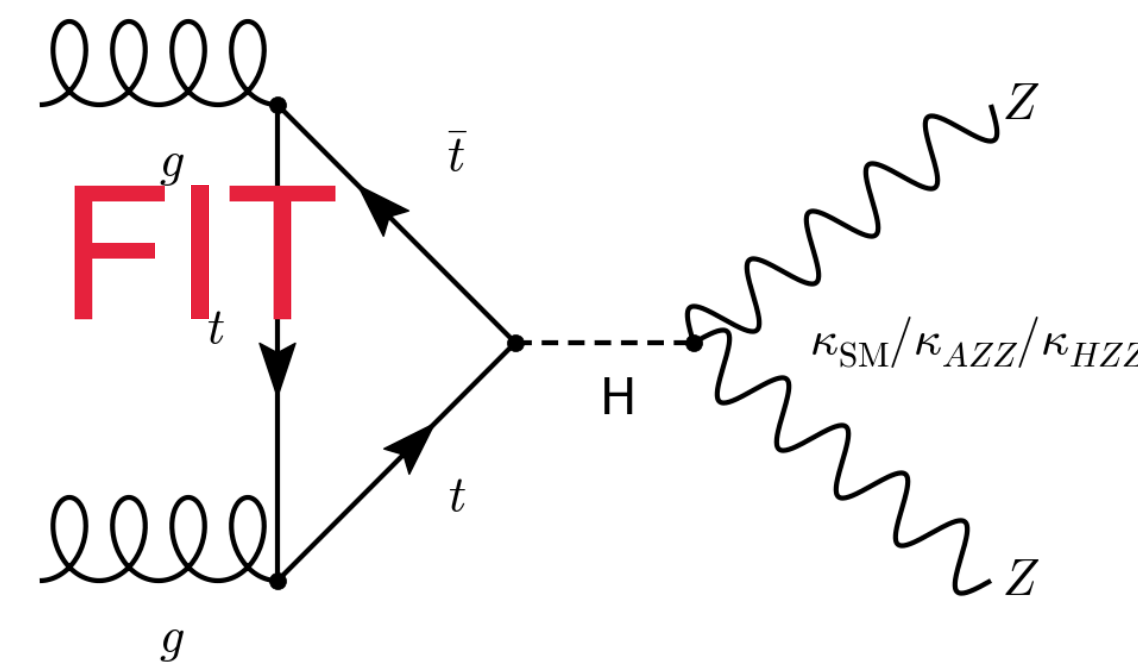
κ_{SM} can contribute to absorb the asymmetry

$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{HVV}, \kappa_{AVV}$ SIMULTANEOUS FIT



Fit configuration	Best-fit $\hat{\kappa}_{HVV}$	Best-fit $\hat{\kappa}_{AVV}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
$\kappa_{Hgg} = 1, \kappa_{SM} = 1$	2.9	± 0.5	-	1.9σ
$\kappa_{Hgg} = 1, \kappa_{SM}$ free	2.1	± 0.3	1.7	1.2σ

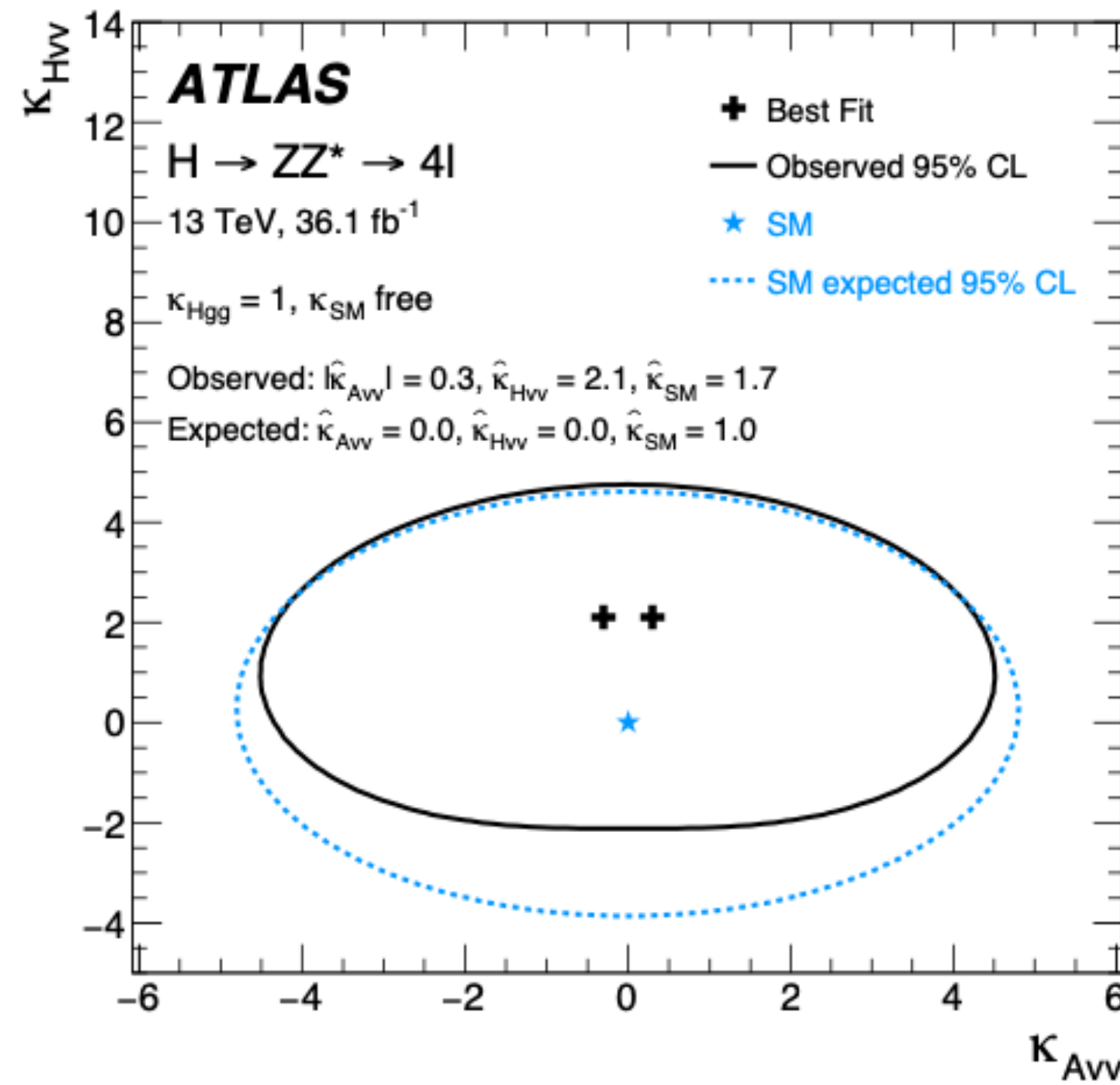
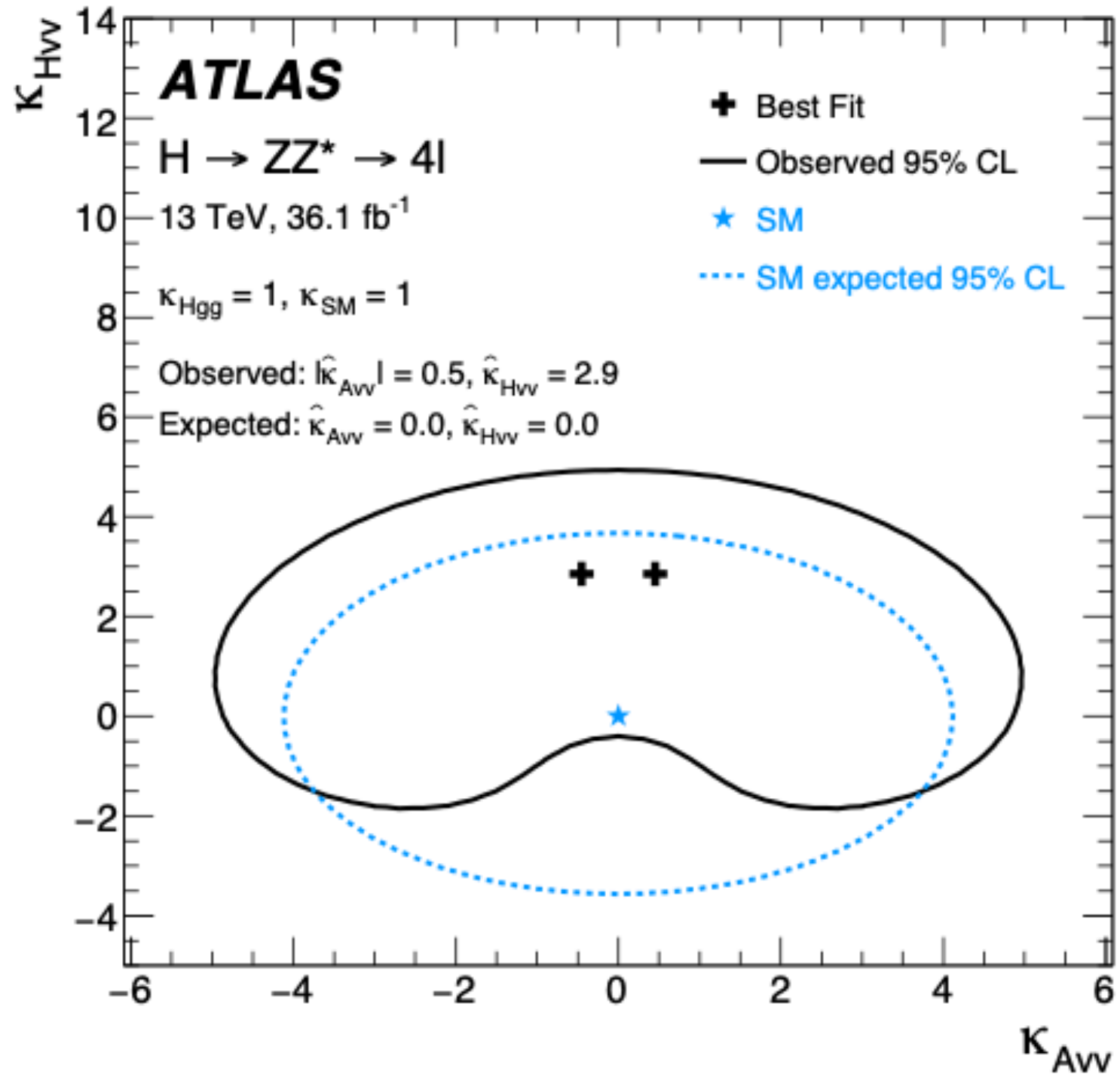
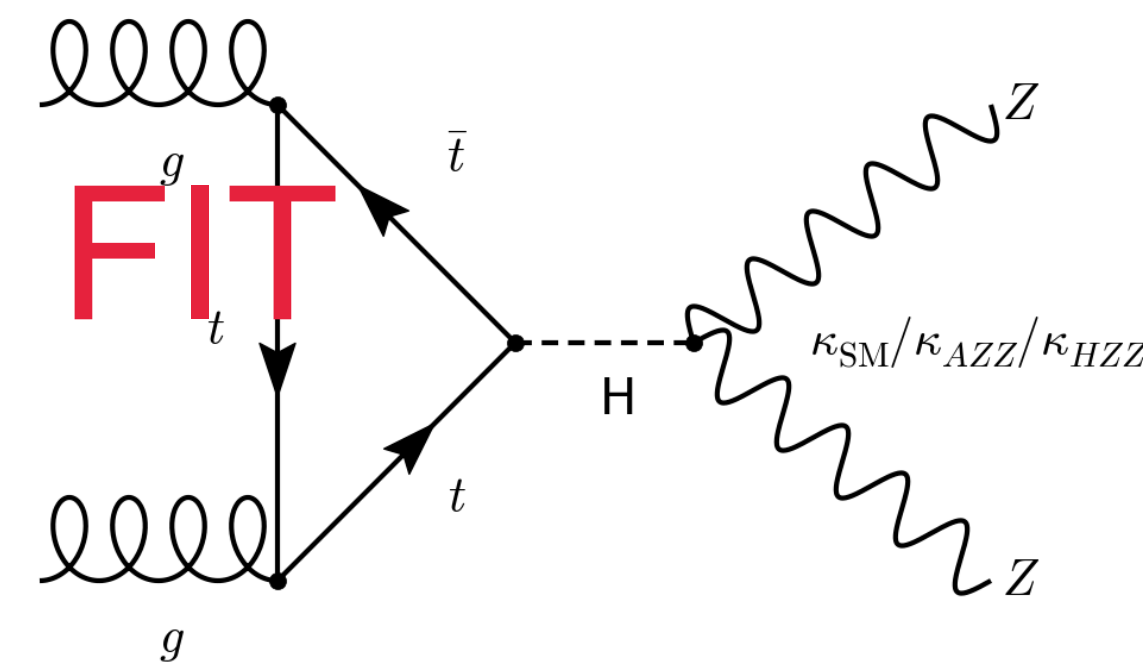
$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{HVV}, \kappa_{AVV}$ SIMULTANEOUS FIT



Fit configuration	Best-fit $\hat{\kappa}_{HVV}$	Best-fit $\hat{\kappa}_{AVV}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
$\kappa_{Hgg} = 1, \kappa_{SM} = 1$	2.9	0.5		1.9σ
$\kappa_{Hgg} = 1, \kappa_{SM}$ free	2.1	± 0.3	1.7	1.2σ

$\hat{\kappa}_{HVV}$ obtained from the 2D scan is similar to the one obtained in the 1D scan.

$H \rightarrow ZZ^* \rightarrow 4l, \kappa_{HVV}, \kappa_{AVV}$ SIMULTANEOUS FIT



Fit configuration	Best-fit $\hat{\kappa}_{HVV}$	Best-fit $\hat{\kappa}_{AVV}$	Best-fit $\hat{\kappa}_{SM}$	Deviation from SM
$\kappa_{Hgg} = 1, \kappa_{SM} = 1$	2.9	± 0.5	1.7	1.9σ
$\kappa_{Hgg} = 1, \kappa_{SM}$ free	2.1	± 0.3	1.7	1.2σ

$\hat{\kappa}_{AVV}$ obtained from the 2D scan is closer to the SM than the corresponding value from the 1D scan.