# Observation of the $\Lambda_b^0 \to J/\psi \Xi^- K^+$ decay

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#### LHCB 2015 *Phys. Rev. Lett.* 115 (2015) 072001 1544 citations!



## Introduction

b hadron decays with charmonium and a baryon allow searching for pentaquarks in  $\psi$ +baryon system in the intermediate resonance structure

LHCb, **2015**: studied  $J/\psi p$  mass from  $\Lambda_b^0 \rightarrow J/\psi p K^-$ 

(full 6D angular analysis with interference between resonances)

# Observed P<sub>c</sub>(4450)<sup>+</sup> and P<sub>c</sub>(4380)<sup>+</sup> pentaguar

pentaquark candidates!

Confirmed later with a <u>model-independent analysis</u> (2016) <u>Also seen</u> in CS  $\Lambda_b^0 \rightarrow J/\psi p \pi^-$  decay (2016)

- **2019:** adding Run-2 data, **9x**  $\Lambda_b^0$  yield. From 1D fit of J/ $\psi$ p mass distribution, 4450 peak is now split into two;

+ observe a new resonance, P<sub>c</sub>(4312)<sup>+</sup>

"Too much data" for a full 6D angular resonance analysis to converge!

#### Introduction





In addition to  $J/\psi p$  system, also the  $J/\psi \Lambda$  system was investigated.

**2020**: 6D full angular analysis by LHCb of  $\Xi_b^- \rightarrow J/\psi \Lambda K^-$  decay <u>revealed evidence</u> for hidden-charm strange pentaquark P<sub>cs</sub>(4459)<sup>o</sup>

<u>CMS-BPH-18-005</u>, <u>JHEP 12 (2019) 100</u>: Based on Run-1, CMS studied the  $B^- \rightarrow J/\psi \Lambda p^-$  decay, <u>data is</u> <u>consistent with no pentaquarks</u> in  $J/\psi \Lambda$  or  $J/\psi p$ 

LHCb 2022: with 6D amplitude analysis of  $B^- \rightarrow J/\psi \Lambda p^-$  decay, observe new strange pentaquark  $P_{cs}(4338)^0 \rightarrow J/\psi \Lambda$  no significant states decaying to  $J/\psi p$ 

It is interesting to note that  $J/\psi \Lambda$  pentaquarks are found to be generally **narrower** than  $J/\psi p$  states (7-17 vs ~10-200 MeV). Even narrower pentaquarks are expected for <u>doubly-strange</u> hidden-charm  $P_{css}$ . Such states can decay into e.g.  $J/\psi \equiv$ 

This motivates our search for decays having  $J/\psi\Xi^-$  in the decay products, i.e.  $\Lambda_b^0 \to J/\psi\Xi^-K^+$ 

#### Data and event selection

Mass constraints applied on  $J/\psi\!\to\mu^+\mu^-,\,\Lambda\!\to\!p\pi^-$  and  $\Xi^-\!\to\!\Lambda\pi^-$ 

 $\Lambda_b^0$  obtained from vertex fit of  $\mu^+\mu^-\Xi^-K^+$ 

Normalization channel is chosen according to the similar decay topology, to reduce the systematic uncertainties associated with the track reconstruction:

 $\Lambda_b^0 \rightarrow \psi(2S)\Lambda$ , with vertex fit of  $\mu^+\mu^-\Lambda\pi^+\pi^+$ , and a requirement on J/ $\psi\pi^+\pi^-$  mass to be close to  $M^{PDG}(\psi(2S))$ 

 $\Lambda_b^0$  vertex should be away from PV in transverse plane

PV selected by smallest angle between  $\Lambda_b^0$  momentum and the line joining PV and  $\Lambda_b^0$  decay vertex

 $\Lambda^0_b$  baryon momentum should be aligned with that line



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#### Optimization of selection criteria

Punzi formula is used for optimization,

with SC recommendation

as it does not rely on **S** normalization

 $f = S/(\frac{463}{13} + 4\sqrt{B} + 5\sqrt{25 + 8\sqrt{B} + 4B})$ 

**S** is number of signal events from MC (double-Gaussian function with common mean)

**B** is expected number of background events in the signal region

Extracted from data with  $m_{PDG}(\Lambda_b^0)\pm 2\sigma_{eff}$  region excluded from the (bkg-only, exponential) fit.

Wrong-sign events are added to the sample to improve statistics.

CS and WS distributions are found to be consistent.

The bkg integral in the signal region is taken as B

<u>Variables</u> Mass windows:

 $m(\Lambda), m(\Xi^{-})$ 

Distance significance between vertices

 $L_{xy}/\sigma_{L_{xy}}(\Xi^-, \Lambda_b^0)$ ,  $L_{xy}/\sigma_{L_{xy}}(\Lambda, \Xi^-)$ ,  $L_{xy}/\sigma_{L_{xy}}(\Lambda_b^0, \mathsf{PV})$ 

Angle between particle momentum and the line passing joining its birth vertex and decay vertex

$$\begin{array}{c} \cos(\overrightarrow{L_{xy}},\overrightarrow{p_T}) \ (\Xi^-,\Lambda_b), \quad \cos(\overrightarrow{L_{xy}},\overrightarrow{p_T}) \ (\Lambda,\Xi^-) \ ,\\ \cos(\overrightarrow{L_{xy}},\overrightarrow{p_T}) \ (\Lambda_b,\mathsf{PV}) \end{array}$$

Transverse momentum

 $p_T(\Lambda_b^0), \ p_T(J/\psi), p_T(\Xi^-), p_T(\Lambda), \ p_T(K^+), \ p_T(\pi^-)$ 

Vertex fit probabilities

 $P_{vtx}(\Lambda_b^0) = P_{vtx}(\Xi^-) = P_{vtx}(\Lambda)$ 

Track impact parameter w.r.t. PV IPS( $\pi$ ), IPS( $K^+$ )

### Calculation of branching fraction ratio



$$\mathcal{B}(\psi(2S) \to J/\psi \pi \pi) = (34.68 \pm 0.30)\%$$
  
 $\mathcal{B}(\Xi \to \Lambda \pi) = (99.887 \pm 0.035)\%$ 

$$\frac{\epsilon_{\psi(2S)\Lambda}}{\epsilon_{J/\psi\Xi^-K^+}} = \frac{4.00 \pm 0.10}{0.79 \pm 0.04} = 5.06 \pm 0.29$$

#### Invariant mass distributions arXiv:2401.16303



 $J/\psi \Xi^{-}K^{+}$  Intermediate invariant mass distributions



8

#### Systematic uncertainties

Source	Uncertainty (%)			
Tracking efficiency	2.3			
$p_{\rm T}(\Lambda_{\rm b}^0)$ spectrum	4.7 $\int$ Different $p_T$ spectra			
Signal model	3.9 $(an the fit model deviation in R = suft upon$			
Background model	$6.7$ $\int$ outy the in model, deviation in k – syst. onc.			
Non- $\psi(2S)$ contribution	2.5			
Limited size of MC samples	5.6			
Selection efficiency	$14.3  ightarrow  ext{Potentially poorly modeled regions}$ of phase space			
Total	18.2			

Total uncertainty is calculated as sum in quadrature of individual sources.

#### Summary

- First observation of  $\Lambda_b^0 \rightarrow J/\psi \Xi^- K^+$ 
  - The first decay to have  $J/\psi \Xi^-$  system in products
- No significant narrow peaks in  $J/\psi\Xi^-$  mass distribution
  - With 46 signal events, our sensitivity is very limited
- Measured branching fraction ratio:

$$\mathcal{R} \equiv \frac{\mathcal{B}(\Lambda_b^0 \to J/\psi \Xi^- K^+)}{\mathcal{B}(\Lambda_b^0 \to \psi(2S)\Lambda)} = [3.38 \pm 1.02 \,(\text{stat}) \pm 0.61 \,(\text{syst}) \pm 0.03 \,(\mathcal{B})]\%$$

$$\frac{\text{arXiv:2401.16303}}{\text{arXiv:2401.16303}}$$

~ same order of magnitude as  $\Lambda_b^0 \rightarrow J/\psi \Lambda \phi$  decay that has similar Feynman diagram:

$$\frac{\mathcal{B}(\Lambda_{\rm b}^0 \to J/\psi \Lambda \phi)}{\mathcal{B}(\Lambda_{\rm b}^0 \to \psi(2{\rm S})\Lambda)} = (8.26 \pm 0.90\,({\rm stat}) \pm 0.68\,({\rm syst}) \pm 0.11(\mathcal{B})) \times 10^{-2}$$

10

# The end.

11

### BACKUP

#### The CMS detector

The central element of the CMS is a **superconducting solenoid** with an internal diameter of 6 m, providing a magnetic field of 3.8 T. Inside the solenoid are **silicon pixel** and **strip detectors**, **electromagnetic** and **scintillation calorimeters**.

Muons are measured using the following detectors: drift tubes, cathode strip chambers with resistive plates.

**Triggers** have 2 levels of information dropout:

- first-level trigger (L1) is a hardware system of triggers that decreases frequency of events to record from 40 MHz to 100 kHz
- high-level trigger (HLT) uses rapid algorithms of event partial reconstruction with decreasing the frequency to 1 kHz



Figure 2: CMS scheme

#### $J/\psi \Xi^- K^+$ invariant mass distribution



Student-T function for signal Exponential for background

arXiv:2401.16303 14

#### Optimization of selection criteria

**Punzi formula** is used for optimization, <u>with SC recommendation</u> as it does not rely on **S** normalization

$$f = S/(\frac{463}{13} + 4\sqrt{B} + 5\sqrt{25 + 8\sqrt{B} + 4B})$$

**S** is number of signal events from MC (double-Gaussian function with common mean)

**B** is expected number of background events in the signal region Extracted from data with  $m_{PDG}(\Lambda_b^0) \pm 2\sigma_{eff}$  region excluded from the (bkg-only, exponential) fit.

Wrong-sign events are added to the sample to improve statistics. CS and WS distributions are found to be consistent.

The bkg integral in the signal region is taken as B

## Optimization of selection criteria for $J/\psi \Xi^- K^+$

- Series of scans over variables performed to find optimal cut values to maximize the expected significance of the signal
- ✓ In each scan, the cut value when f takes the largest value is recorded and used in the following scans
- ✓ When iteration shows the same result (cut values) as the previous one, the optimization is complete
- ✓ Selection criteria for normalization channel are chosen similar (as close as possible) to those found for the signal channel

<u>Variables</u>

Mass windows:

m(Λ), m(Ξ<sup>-</sup>)

Distance significance between vertices

 $L_{xy}/\sigma_{L_{xy}}(\Xi^-, \Lambda_b^0)$ ,  $L_{xy}/\sigma_{L_{xy}}(\Lambda, \Xi^-)$ ,  $L_{xy}/\sigma_{L_{xy}}(\Lambda_b^0, \mathsf{PV})$ 

Angle between particle momentum and the line passing joining its birth vertex and decay vertex

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Transverse momentum

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Vertex fit probabilities

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Track impact parameter w.r.t. PV IPS( $\pi$ ), IPS( $K^+$ )

### Systematic uncertainties

- 1) Uncertainty of efficiency ratio due to limited MC statistics
- 2) Signal model choice: try several alternative models, take the largest variation in R as systematics
  - Student-T is baseline, alternatives are
    - o Double-gaussian
    - o Johnson PDF
- 3) Background model choice: several alternative models  $\rightarrow$  largest variation in R
  - Exp is baseline, alternatives are
    - o 2<sup>nd</sup> degree polynomial
    - Modified threshold pdf  $(x-x^0)^{\alpha} \cdot exp$
    - Modified threshold pdf  $(x-x^0)^{\alpha} \bullet Pol_1$
- 4) Tracking efficiency:

the pT spectra of the harder of the two tracks are found to differ significantly between signal and norm. channels → **conservatively** taking 2.3% as additional systematic as if there were different number of tracks in 2 channels



Systematic uncertainties - Potential non-psi(2S) contribution



To estimate background under  $\psi(2S)$  we use *sPlot* method to subtract the background under  $\Lambda_b^0$ . The m(J/ $\psi\pi\pi$ ) range was expanded to 5 $\sigma$  around mPDG( $\psi(2S)$ ). Integral of bckg function in baseline region [ $|m(J/\psi\pi\pi) - mPDG(\psi(2S))| < 11.1 \text{ MeV}$ ] is  $30\pm18$ 

Thus, the additional systematic uncertainty is 30/1179 = **2.5% 1179** - the signal yield for R measurement cuts

#### Systematic uncertainties - Selection efficiency

Variable	10% drop (20% drop)	R, %	$\mathcal{R}_{uncor}$ , %	$\sqrt{d^2 - (\delta d)^2}/3.38\%$	
$p_{\mathrm{T}}(\mu)$	4.45 GeV	$3.50 \pm 1.12$	$3.50 \pm 0.53$	-	Change in R:
$p_{\rm T}(\mu)$	(4.8 GeV)	$3.03 \pm 1.06$	$3.03\pm0.42$	-	d = 2.68 - 3.38 = 0.70%
$p_{\rm T}({\rm J}/\psi)$	10.5 GeV	$3.44 \pm 1.14$	$3.44\pm0.32$	-	
$p_{\rm T}({\rm J}/\psi)$	(12.0 GeV)	$2.68 \pm 1.14$	$2.68\pm0.52$	14.3%	¥
$P_{vtx}(J/\psi)$	19%	$3.25 \pm 1.07$	$3.25 \pm 0.41$	_	Its uncertainty:
$P_{vtx}(J/\psi)$	(30%)	$3.35 \pm 1.14$	$3.35 \pm 0.56$	_	$\delta d = 0.52\%$
$IPS(K^+ \Lambda_{b}^{0})$	2.8	$3.30 \pm 1.04$	$3.30 \pm 0.11$	_	0.0270
$IPS(K^+ \Lambda_b^0)$	(3.45)	$3.84 \pm 1.20$	$3.84\pm0.67$	_	
$p_{\rm T}(\pi_{\Xi}^-)$	0.55 GeV	$3.60 \pm 1.13$	$3.60\pm0.45$	_	· · · · · · · · · · · · · · · · · · ·
$p_{\rm T}(\pi_{\Xi}^{\Xi})$	(0.67 GeV)	$3.23 \pm 1.15$	$3.23 \pm 0.43$	_	Square root difference
$cos(\overrightarrow{L_{xy}}, \overrightarrow{p_{\rm T}})({\rm J}/\psi_{\rm -}PV)$	0.9975	$3.40 \pm 1.07$	$3.40 \pm 0.59$	_	between them:
$cos(\overrightarrow{L_{xy}}, \overrightarrow{p_{T}})(J/\psi_{-}PV)$	(0.9985)	$3.77 \pm 1.27$	$3.77\pm0.50$	_	$\sqrt{d^2 - (\delta d)^2} = 0.47\%$
$L_{xy}/\sigma_{L_{xy}}(J/\psi_P V)$	11.5	$2.95 \pm 1.03$	$2.95\pm0.45$	-	
$L_{xy}/\sigma_{L_{xy}}(J/\psi_P V)$	(16.0)	$2.90\pm1.10$	$2.90\pm0.53$	-	$\downarrow$
Baseline		$3.38 \pm 1.02$	3.38		Additional systematic

We strengthen the cut and evaluate the uncertainty in the phase space where the signal events are located. We vary the each cut individually, strengthening the requirement until the efficiency is at 80% with respect to the nominal value and at 90% as a cross-check.

19

uncertainty:

0.47/3.38 = 14.3%