



#### **Cosmological Constant Suppression in Non-Stationary Scalar Covariant State**

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- Introduction
- Coherent spatial-temporal structure of fluctuations
- Coherent spatial-temporal structure in finite volume with scale cut-off
- Numerical estimates
- Conclusion

- Cosmological constant problem: according to the data  $10^{-3}$  eV, based on estimates in theoretical models  $10^{18}$  GeV;
- No theoretical models with vacuum energy density suppression ⇒ Zeldovich: the cosmological constant ↔ density of vacuum energy generated by ZPM;
- 4D isotopic model in contrast to 3D isotropic

#### Zero-point modes of scalar field and spatial covariance

• For SE tensor in ultra-relativistic expressions ( $\Lambda_P \gg m$ ):

$$\langle 0 | T_{\alpha\beta} | 0 \rangle \approx \frac{1}{3} \,\delta_{\alpha\beta} \int_{0}^{\Lambda_{P}} \frac{4\pi p^{2} dp}{(2\pi)^{3}} \,\frac{p^{2}}{2p} \Big( 1 - \frac{m^{2}}{2p^{2}} \Big) = \frac{1}{3} \,\delta_{\alpha\beta} \cdot \frac{\Lambda_{P}^{4}}{16\pi^{2}} \Big( 1 - \frac{m^{2}}{\Lambda_{P}^{2}} \Big),$$

$$\langle 0 | T_{00} | 0 \rangle \approx \int_{0}^{\Lambda_{P}} \frac{4\pi p^{2} dp}{(2\pi)^{3}} \frac{p}{2} \left( 1 + \frac{m^{2}}{2p^{2}} \right) = \frac{\Lambda_{P}^{4}}{16\pi^{2}} \left( 1 + \frac{m^{2}}{\Lambda_{P}^{2}} \right).$$

• Therefore, the ratio of isotropic pressure:

$$w = \frac{p}{\rho} \approx \frac{1}{3} \left( 1 - \frac{2}{m^2} \Lambda_P^2 \right) \rightarrow w_{\text{rad}} = \frac{1}{3} \text{ at } \frac{m}{\Lambda_P} \rightarrow 0.$$

• Analogously, in the non-relativistic limit:

$$w = \frac{p}{\rho} \approx \frac{3}{5} \frac{\Lambda_P^2}{m^2} \left( 1 - \frac{23}{35} \frac{\Lambda_P^2}{m^2} \right) \to w_{\rm m} = 0 \quad \text{at} \frac{\Lambda_P}{m} \to 0.$$

#### Scalar field and spatial-temporal covariance

• Suppose true spatial-temporal structure (Wick rotation, no poles):

$$\left\langle \operatorname{vac} \left| T_{\mu\nu}(x) \right| \operatorname{vac} \right\rangle = \int \frac{\mathrm{d}^4 p_E}{(2\pi)^4} \, \frac{1}{p_E^2 + m^2} \left( p_\mu^E p_\nu^E - \frac{1}{2} \, g_{\mu\nu}^E(p_E^2 + m^2) \right)$$

• Replacement:

$$p^E_\mu p^E_\nu \mapsto \frac{1}{4} g^E_{\mu\nu} p^2_E.$$

• Therefore:

$$\langle \operatorname{vac} | T_{\mu\nu}(x) | \operatorname{vac} \rangle = \frac{1}{4} g_{\mu\nu} \int \frac{\mathrm{d}^4 p_E}{(2\pi)^4} \left( 1 + \frac{m^2}{p_E^2 + m^2} \right) \Rightarrow w_{\operatorname{vac}} = -1$$

• True structure of spatial-temporal quantum fluctuations in the vacuum supposes the covariant four-dimensional formulation.

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### Finite volume and spatial structure

• The scalar field in a finite volume:

$$S = V_{[3]} \int dt \, \frac{1}{2} \left( (\partial_t \phi(t))^2 - m^2 \phi^2(t) \right)$$

• Using oscillator coordinate, mass, frequency and scale of energy:

$$\langle 0 | m^2 \phi^2 | 0 \rangle = \langle 0 | (\partial_t \phi)^2 | 0 \rangle = \frac{1}{2} \frac{m}{V_{[3]}},$$

• So, the energy density equals:

$$\langle 0 | T_{00} | 0 \rangle = \frac{1}{2} \langle 0 | (\partial_t \phi)^2 + m^2 \phi^2 | 0 \rangle = \frac{1}{2} \frac{m}{V_{[3]}},$$

• while the pressure is given by:

$$p = \frac{1}{2} \langle 0 | (\partial_t \phi)^2 - m^2 \phi^2 | 0 \rangle = 0.$$

 Does not correspond to the vacuum: the parameter fits the dust, i.e. particles with zeroth pressure

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#### Covariant model

• Action for the field in the **finite** volume in the **covariant 4d** form:

$$S_{\Lambda} = \frac{1}{\Lambda^3} \int \mathrm{d}\tau \frac{1}{2} \Big( (\partial_{\tau} \phi)^2 \Big\langle \partial_{\mu} \tau \partial_{\nu} \tau \Big\rangle g^{\mu\nu} - m^2 \phi^2 \Big),$$

• the covariant ST-structure in the model is defined by:

$$\left\langle \partial_{\mu} \tau \partial_{\nu} \tau \right\rangle = \frac{1}{4} g_{\mu\nu}, \quad \left\langle (\partial_{\lambda} \tau)^2 \right\rangle = 1$$

Analogously for the basic state with no quanta of oscillator:

$$\langle \operatorname{vac} | T^{\Lambda}_{\mu\nu} | \operatorname{vac} \rangle = g_{\mu\nu} \langle \operatorname{vac} | \left( -\frac{1}{4} \left( \partial_{\tau} \phi \right)^2 + \frac{1}{2} m^2 \phi^2 \right) | \operatorname{vac} \rangle = g_{\mu\nu} \frac{1}{8} m \Lambda^3$$

and the bare value of energy density:

$$\rho_{\rm vac}^{\rm bare} = \langle \operatorname{vac} | T_{00}^{\Lambda} | \operatorname{vac} \rangle = \frac{1}{8} m \Lambda^3 > 0.$$

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### Non-stationary coherent state

 The contribution of vacuum state to the average stress-energy tensor is suppressed as:

$$\langle T^{\Lambda}_{\mu\nu} \rangle_{\rm vac} = \langle \operatorname{vac} | T^{\Lambda}_{\mu\nu} | \operatorname{vac} \rangle \cdot e^{-\langle n \rangle} = \rho_{\rm vac}^{\rm bare} \cdot e^{-\langle n \rangle} = \frac{1}{8} m \Lambda^3 \cdot e^{-\langle n \rangle},$$

• the energy of coherent state in the final volume:

$$\langle E \rangle = \omega_{\rm osc} \left( \langle n \rangle + \frac{1}{2} \right) = m \left( \langle n \rangle + \frac{1}{2} \right),$$

• the same quantity in terms of energy density within a finite volume:

$$\langle E \rangle = V_{[3]} \rho_{\text{vac}}^{\text{bare}} \left( \langle n \rangle + \frac{1}{2} \right) \cdot 2 = V_{[3]} \frac{1}{4} m \Lambda^3 \left( \langle n \rangle + \frac{1}{2} \right),$$

• Therefore, the reference volume:

$$V_{[3]} = \frac{4}{\Lambda^3}$$

 The contribution of vacuum state to the average stress-energy tensor is suppressed as:

$$\rho_{\rm vac} \sim (10^{-3} \,{\rm eV})^4, \qquad \Lambda \sim 10^{16} \,{\rm GeV}\,.$$

• Then  $\langle n \rangle \sim 250 \Rightarrow$  the energy of initial coherent state in the reference volume:

$$\langle E \rangle \sim \langle n \rangle m \sim \tilde{m}_{\rm Pl}$$
.

 Two energy scales the reduced Planckean mass and the scale of inflation plateau.

## Conclusion

- We have shown the bare value of ZP can be suppressed under the following conditions (+covariant formalism):
  - the non-stationary coherent state of scalar field in the expanding universe
  - the finite volume
  - natural estimates for the cut-off scale
- The model involves at least two scales of energy

# Thank you for your attention!